

Magnetosfere di stelle di neutroni in accrescimento e di magnetars

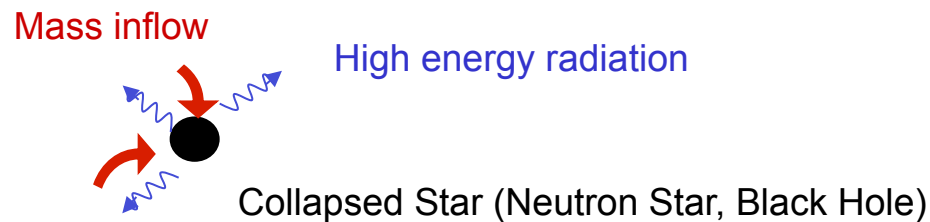
Luigi Stella

INAF – Osservatorio Astronomico di
Roma

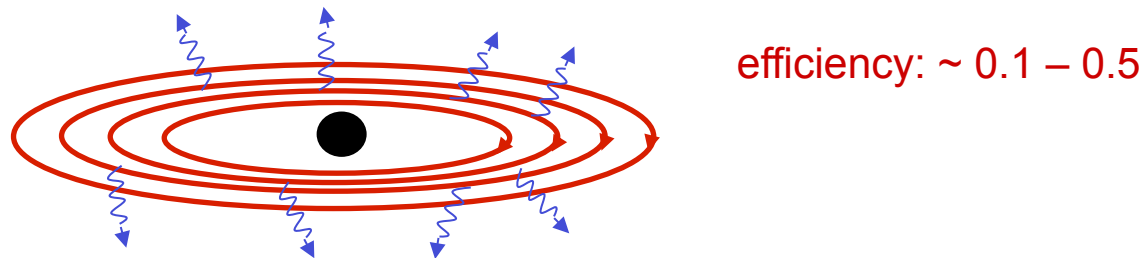
Accretion into collapsed objects

Matter accretion toward collapsed stars:

→ Energy production with high efficiency!



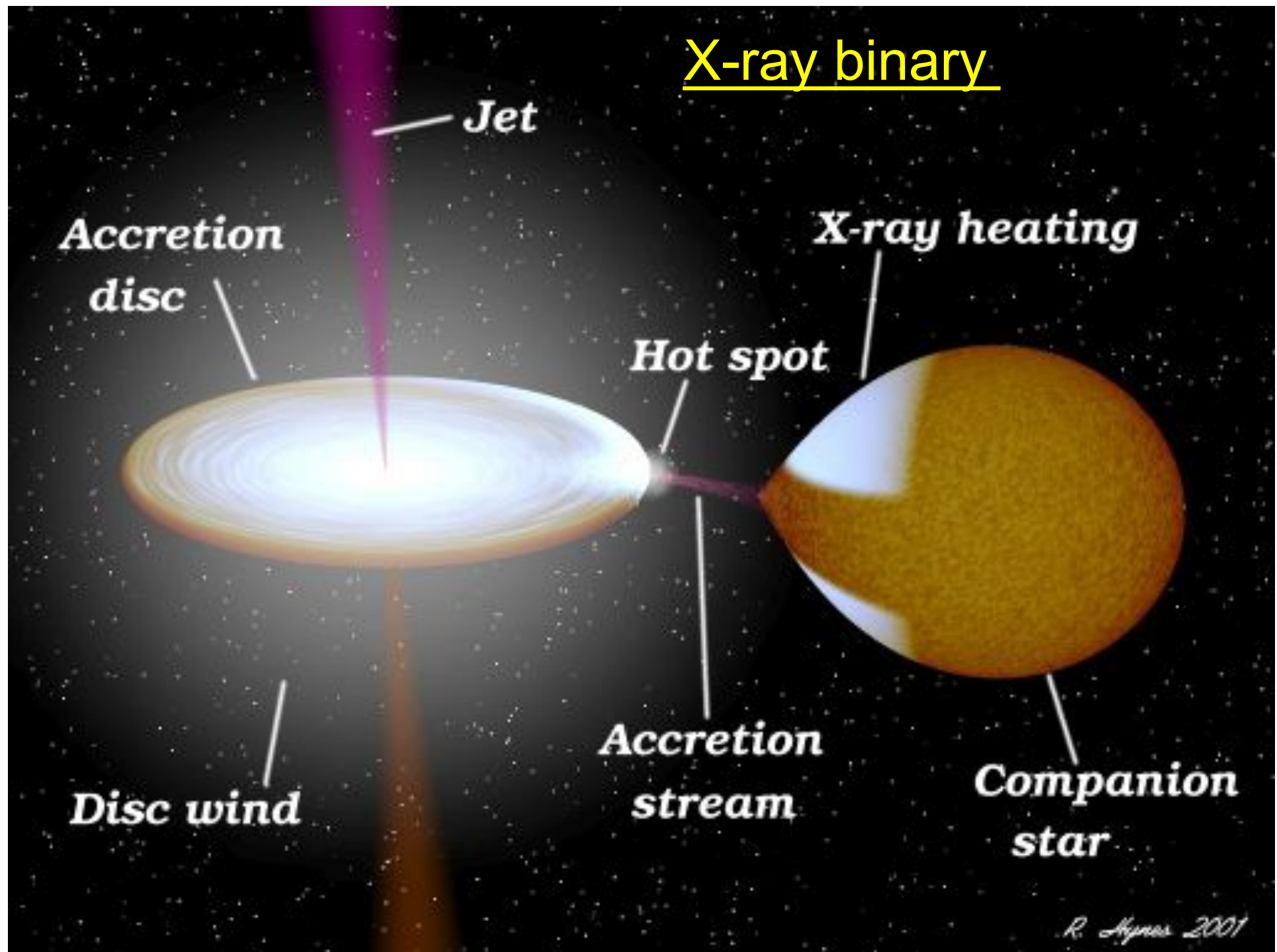
→ If matter has enough L an accretion disk forms:



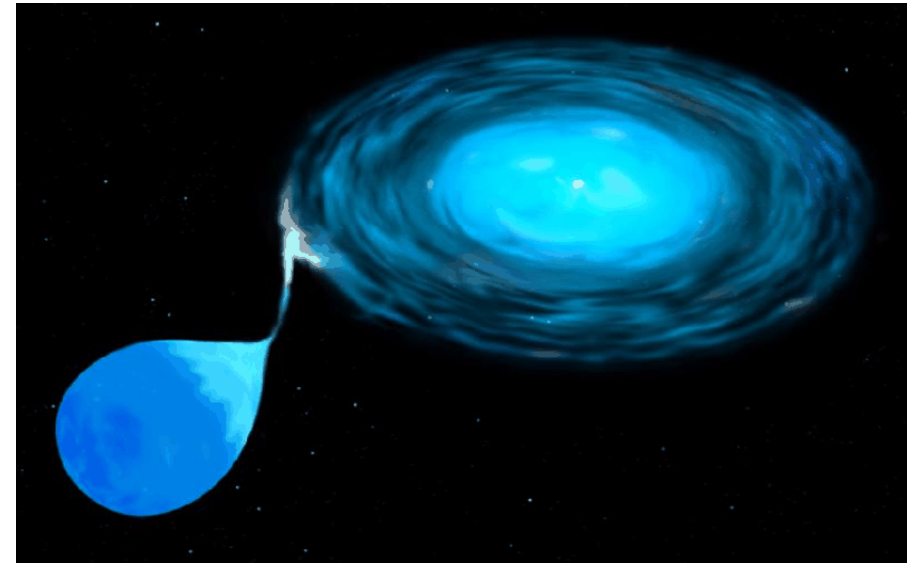
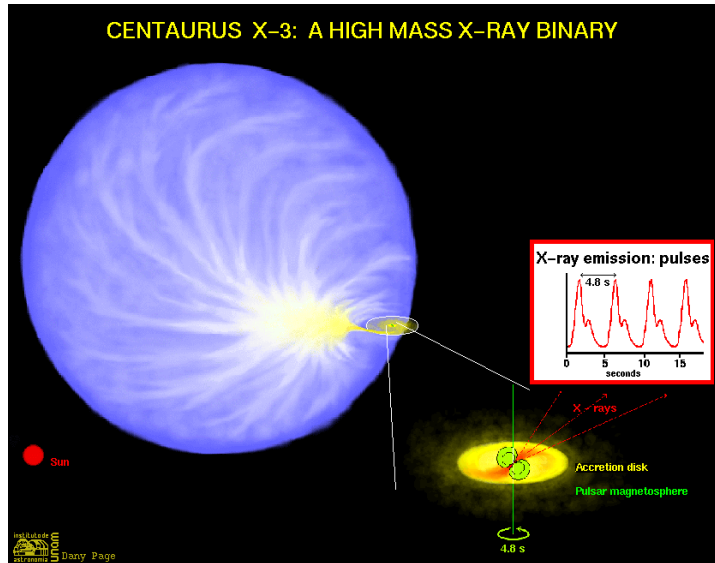
Regions of strong gravitational field accessible to study.

Systems are not “clean”: a variety of effects (plasma, radiative transfer, fluidodynamics, MHD) complicate matters.

X-ray binary



High Mass XRB vs. Low Mass XRB



- HMXB : companion mass $M_2 > M_{\odot}$
- Luminous, young, early type (O, B)
- Stellar wind fed systems
- Often contains an X-ray pulsar

- LMXB : companion mass $M_2 \leq M_{\odot}$
- Faint, old, late spectral type (K, M)
- Roche-Lobe / Disk fed systems
- Often transient system

ALSO

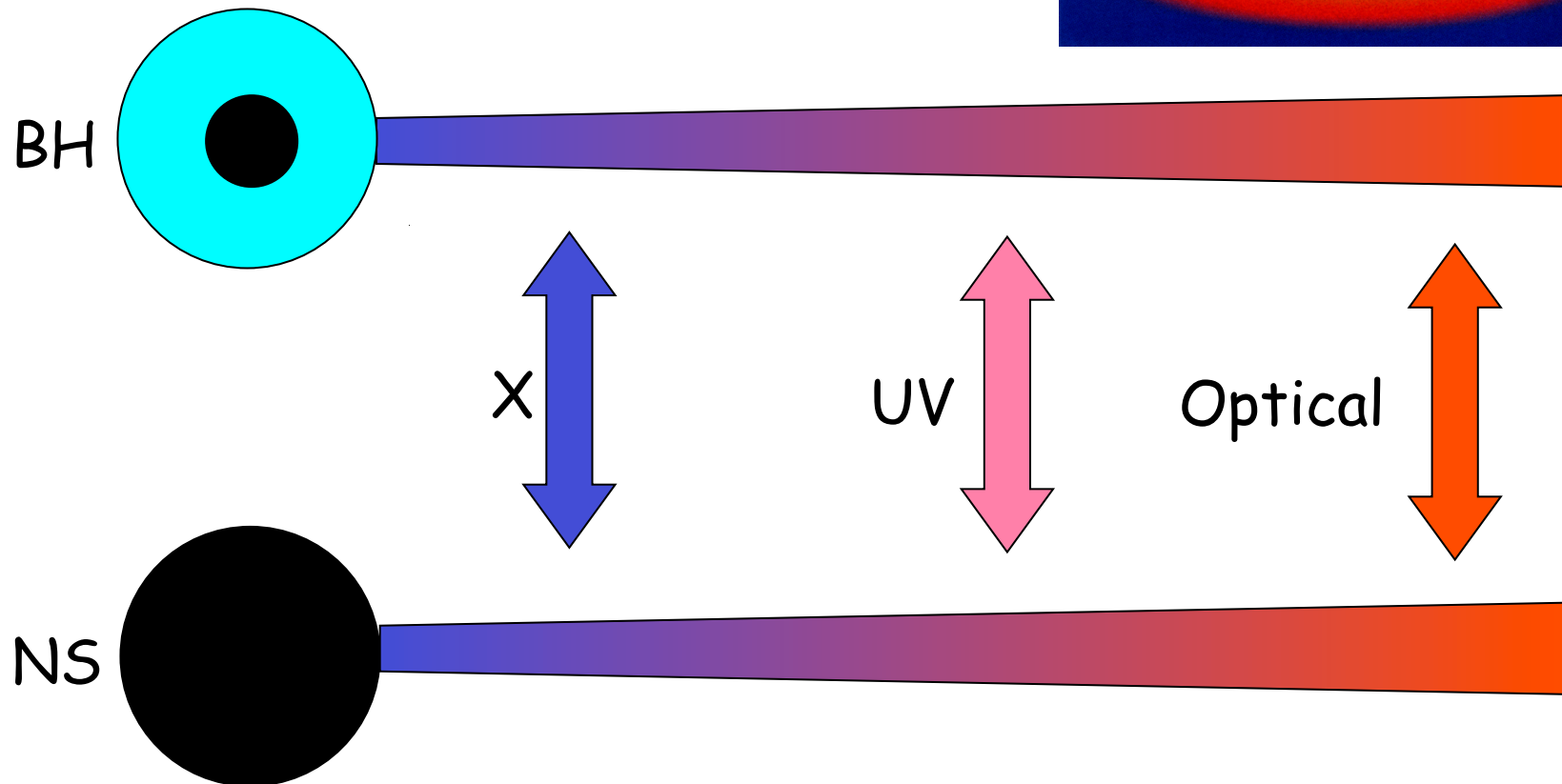
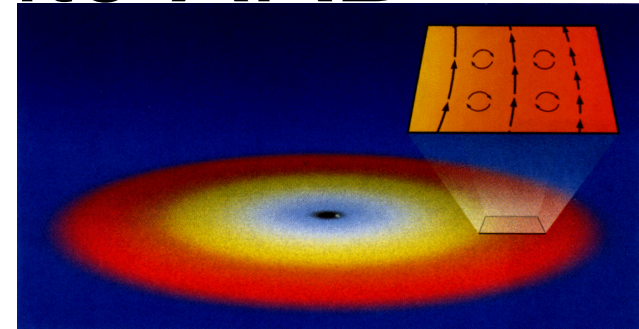
Neutron Star XRB vs. Black Hole XRB

Disk accretion onto XRB

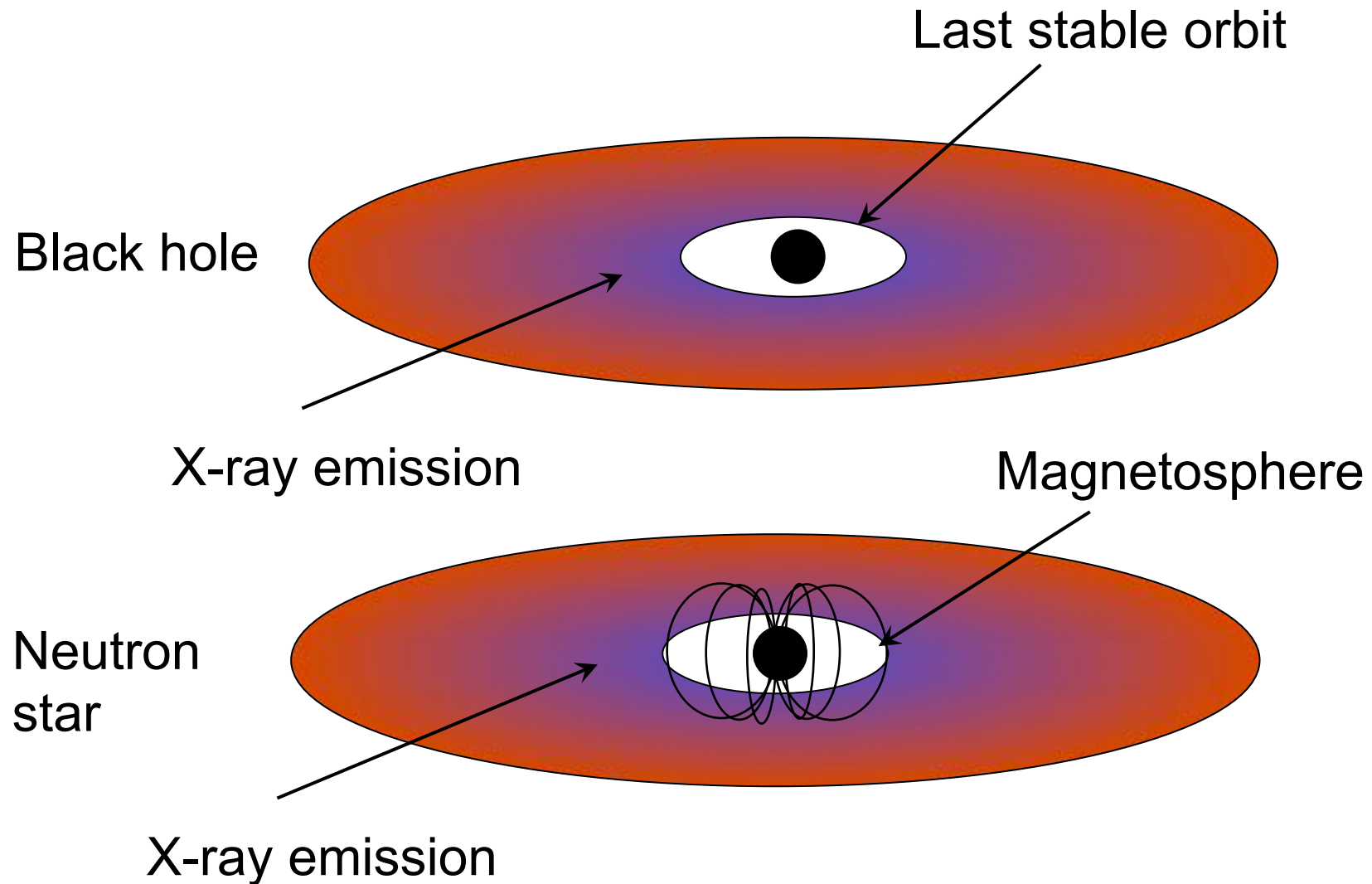
Accretion disk structure

No ordered B field in BH

BH or NS?

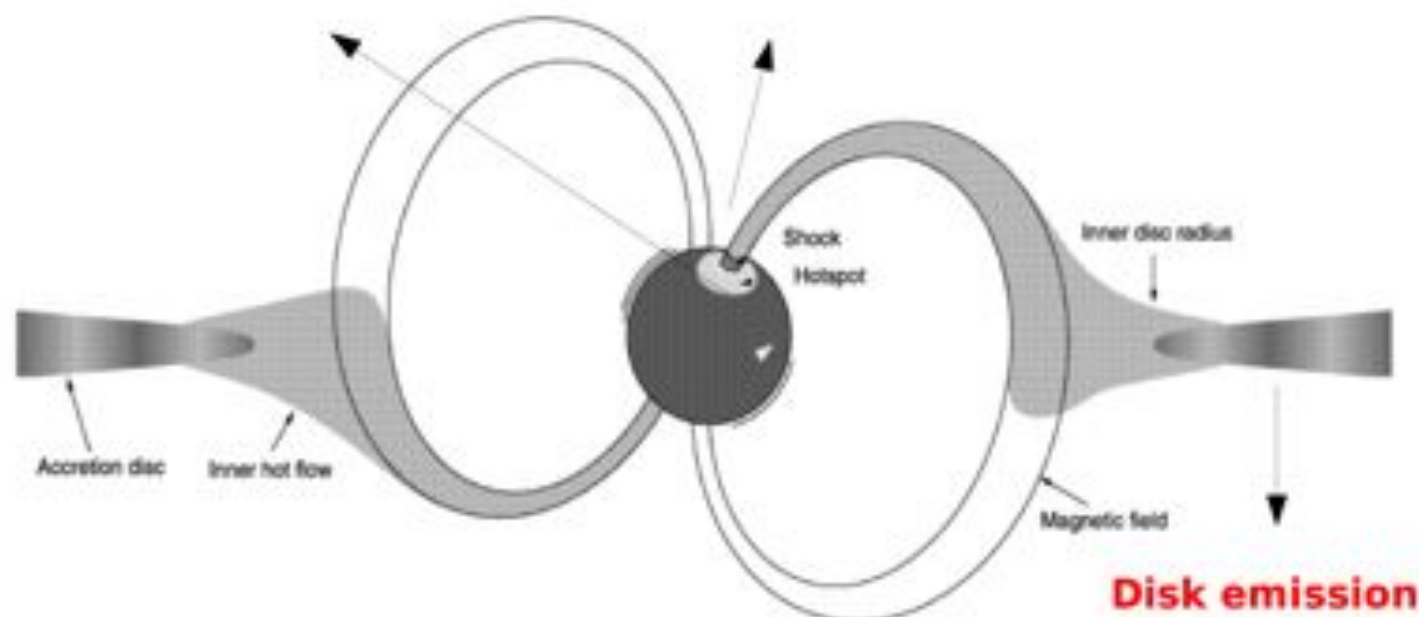


BH vs. NS



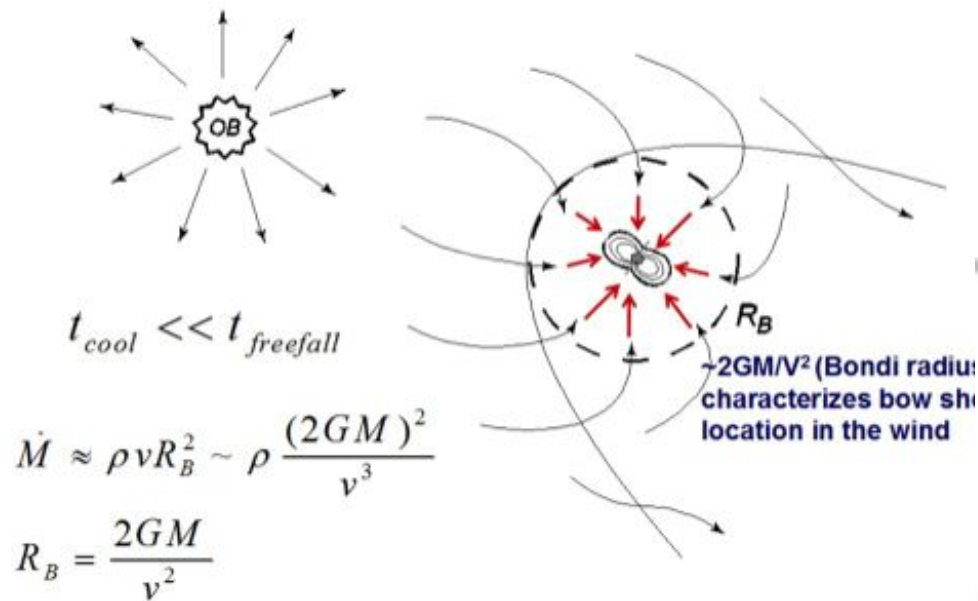
Hot spots – Reprocessing of the Irradiating hard photons into thermal radiation

Accretion Column – Accreting plasma heated above NS surface $\rightarrow$ Comptonization of soft photons from the surface



Spherical accretion model

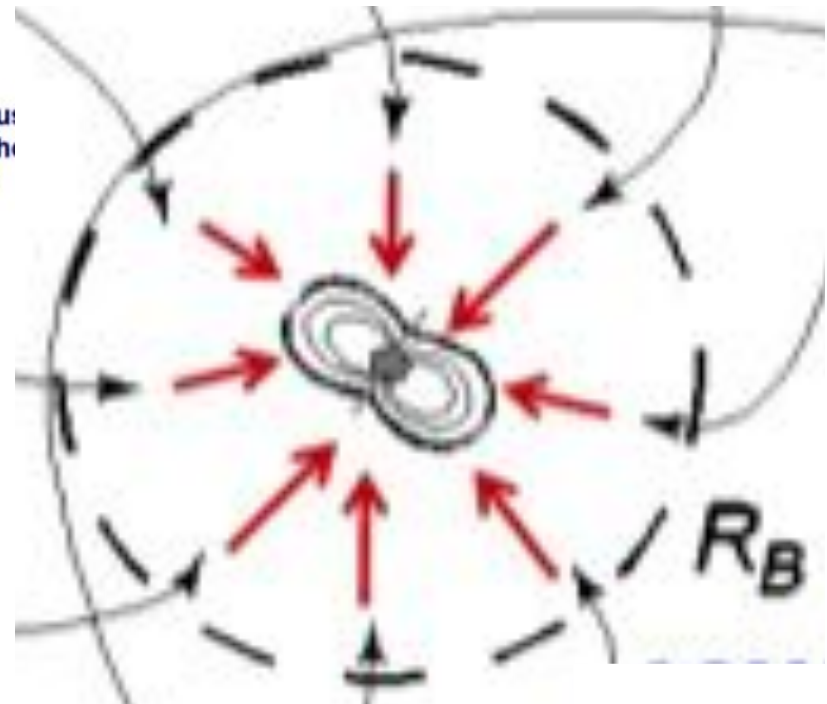
Accretion Bondi-Hoyle-Littleton



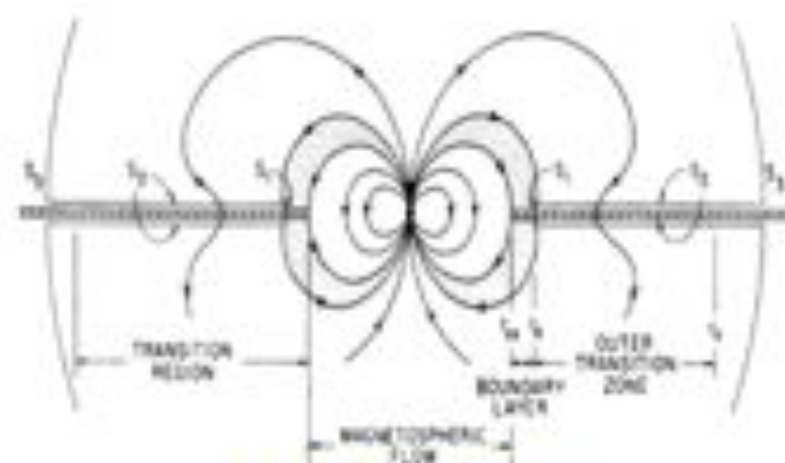
To work out Rm set:

Ram pressure = Magnetic pressure

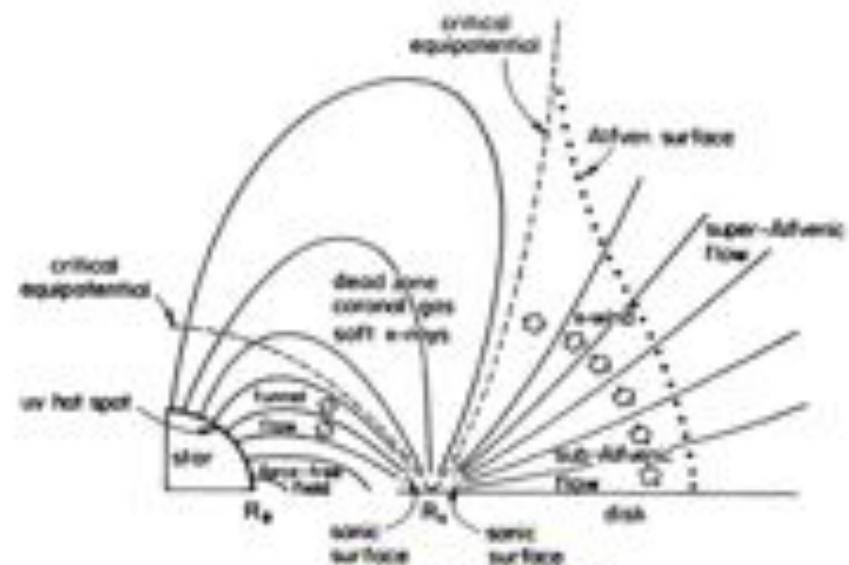
$$r_A = \left(\frac{\mu^4}{2GM\dot{M}^2} \right)^{1/7}$$



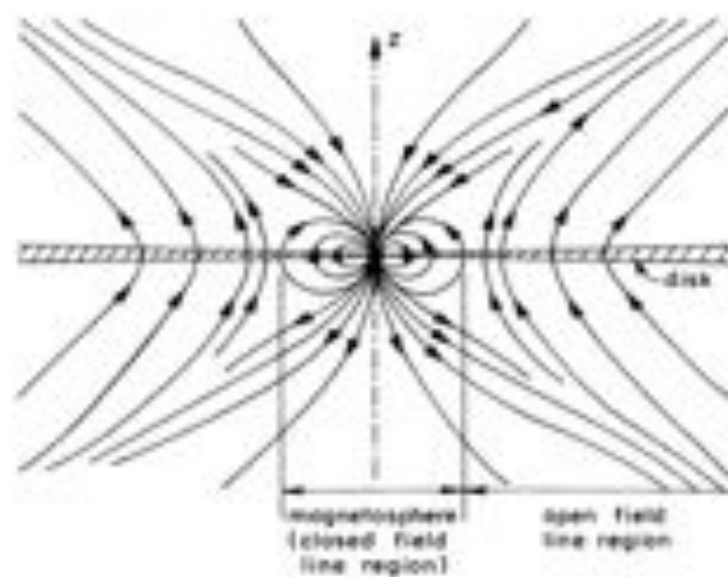
$$r_{cor} = \left(\frac{GM P^2}{4\pi^2} \right)^{1/3}$$



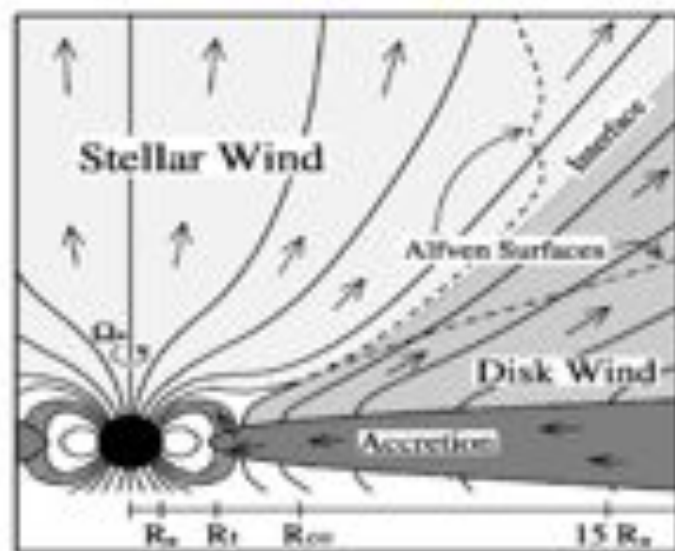
Ghosh & Lamb 1979



Shu et al. 1994



Lovelace et al. 1995



Matt & Pudritz 2005

$$r_{\text{cor}} = \left(\frac{G M P^2}{4 \pi^2} \right)^{1/3}$$

- For disk accretion neutron star, the spin will reach a equilibrium spin period when the co-rotation radius equal the magnetosphere radius.

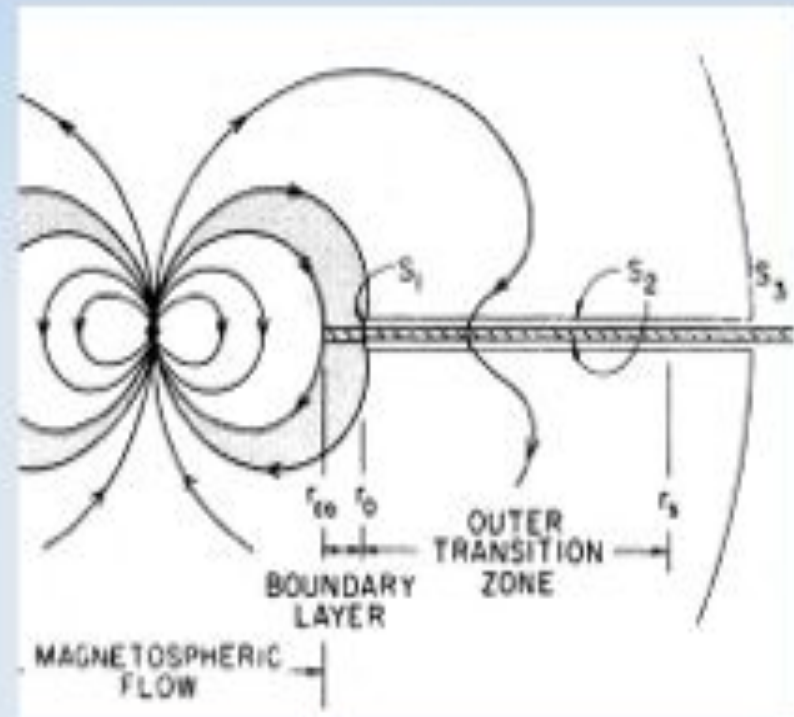
$R_m > R_c$ spin down propeller phase
 $R_m < R_c$ spin up accretion phase

$$R_m = \left[\mu^4 / 2 G M_n M^2 \right]^{1/7}$$

$$R_c = \left[G M_n P^2 / 4 \pi^2 \right]^{1/3}$$

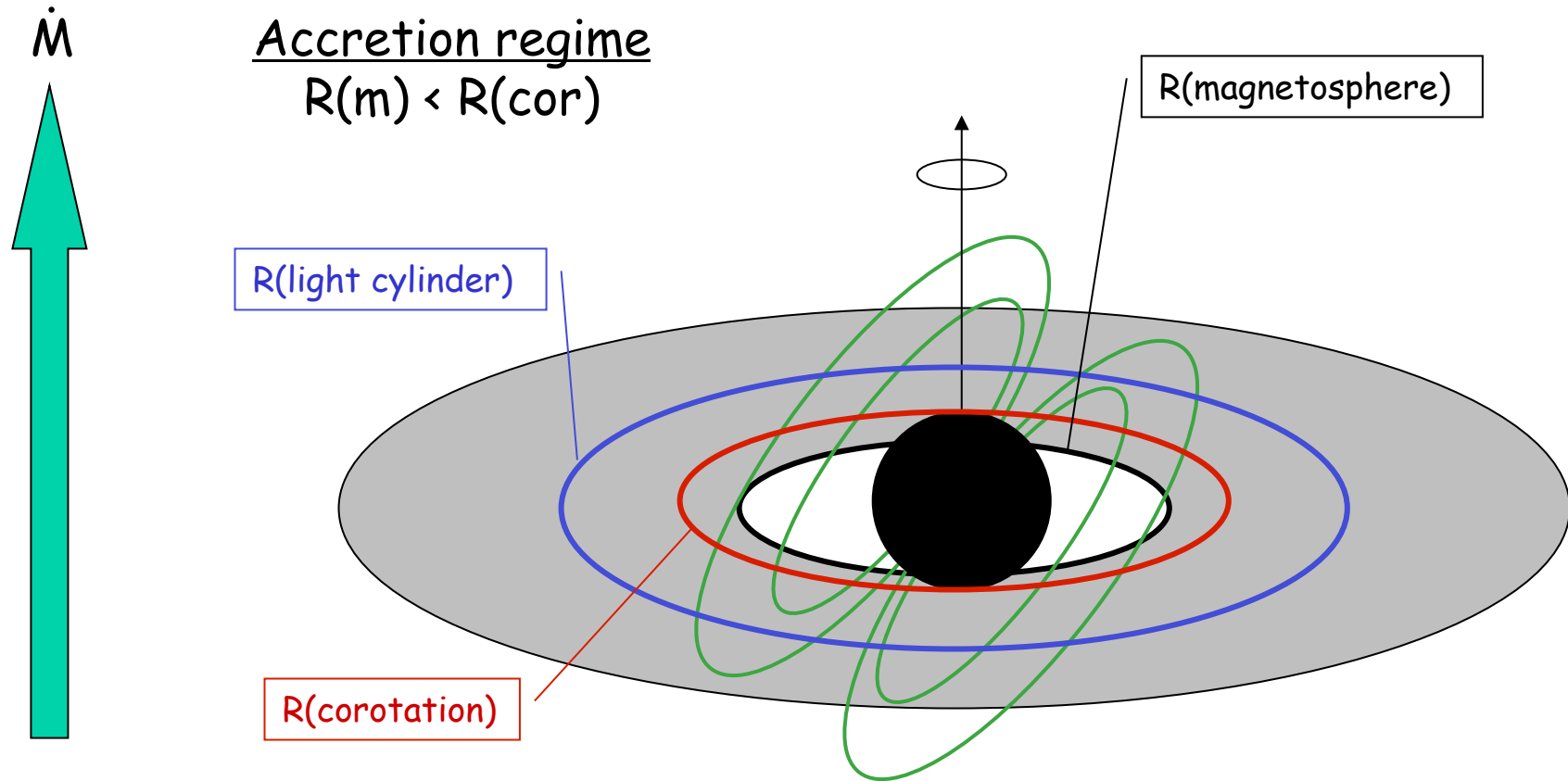
$$R_m \sim R_c$$

$$P_{\text{eq}} = 2 \pi \sqrt{2} \mu^{6/7} (2 G M_n)^{5/7} M^{-3/7}$$



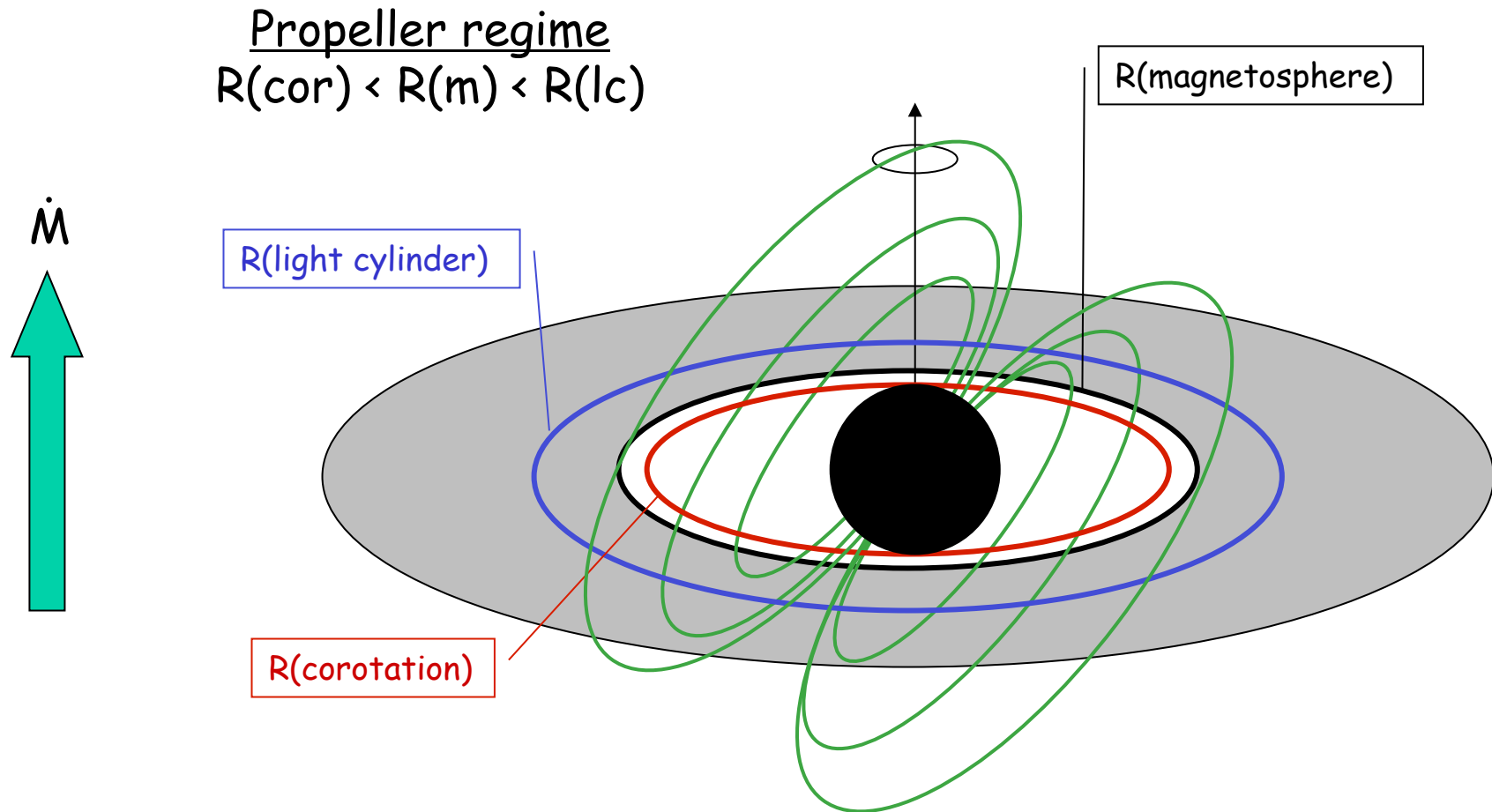
Different regimes for a rotating magnetic NS: 1

(Illarionov & Sunyaev 1975)



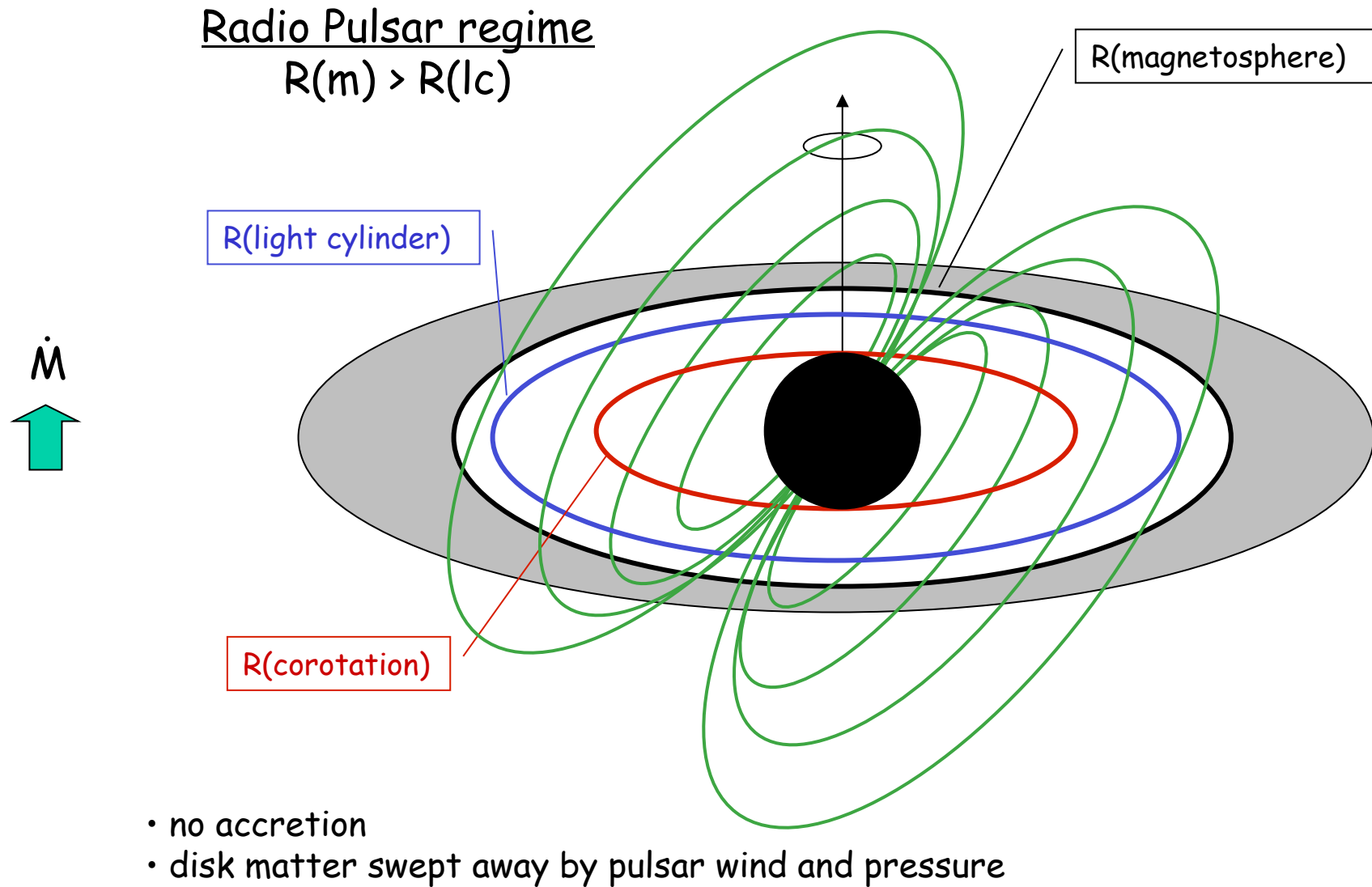
- accretion onto NS surface (magnetic poles)
- energy release $L = GMM(\dot{M})/R^*$

Different regimes for a rotating magnetic NS: 2



- centrifugal barrier closes (B-field drag stronger than gravity)
- matter accumulates or is ejected from $R(\text{m})$
- accretion onto $R(\text{m})$: lower gravitational energy released

Different regimes for a rotating magnetic NS: 3



Why are X-ray transients important ?

Outburst X-ray luminosities: $L_x \sim 10^{37-39} \text{ erg/s}$,

Quiescent X-ray luminosities: $L_x \sim 10^{31-33} \text{ erg/s}$

- Physics of accreting collapsed stars over a very large range of mass inflow rates, which is not accessible to persistent sources

Different regimes in Hard XRTs

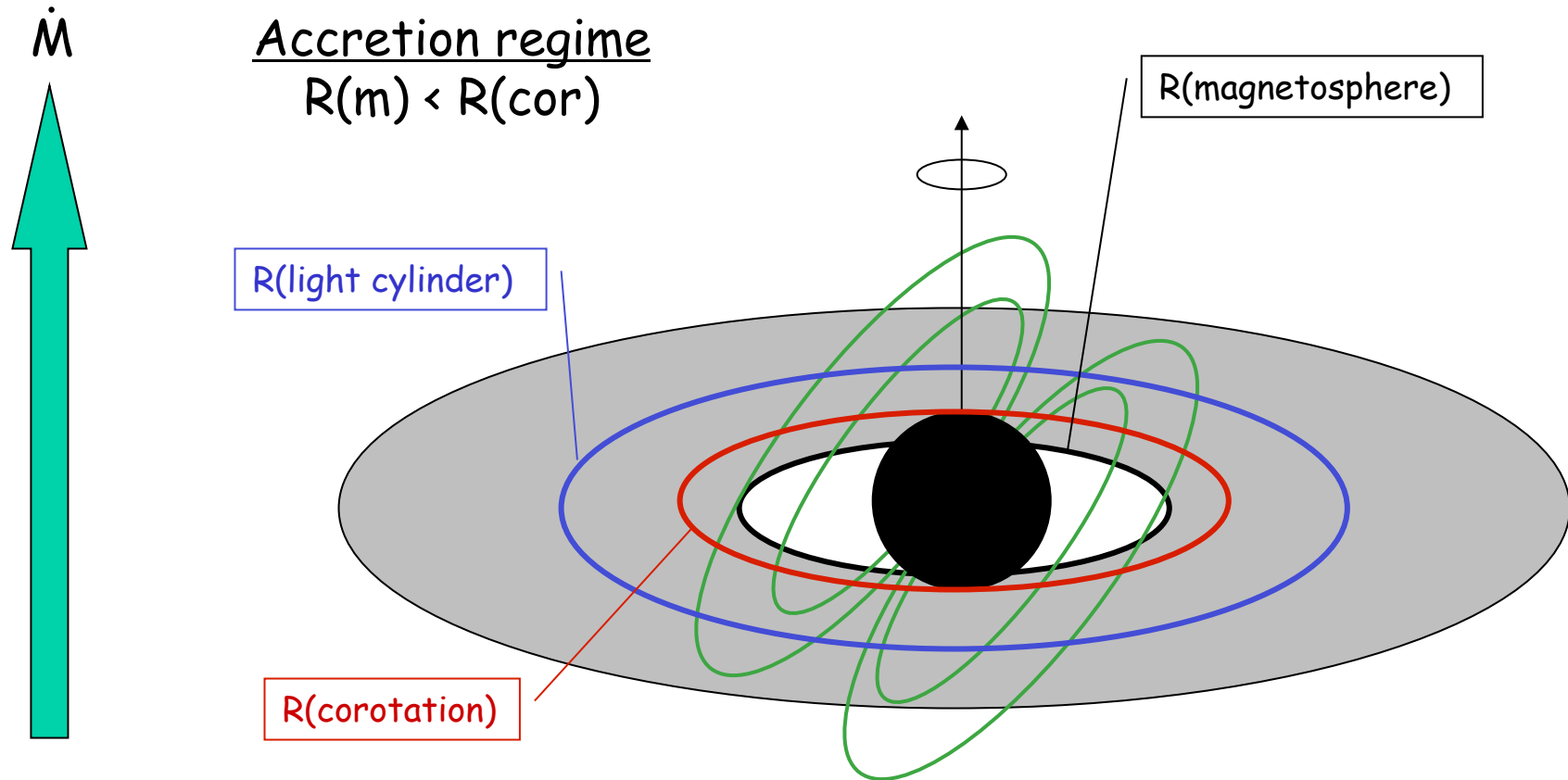
(X-ray pulsar/ Be star binaries)

(Stella, White & Rosner 1986)

- Centrifugal barrier likely closes close to an outburst end: X-ray flux decay should steepen suddenly
- Self-consistency check of interpretation, if $P(\text{spin})$, B and distance are measured

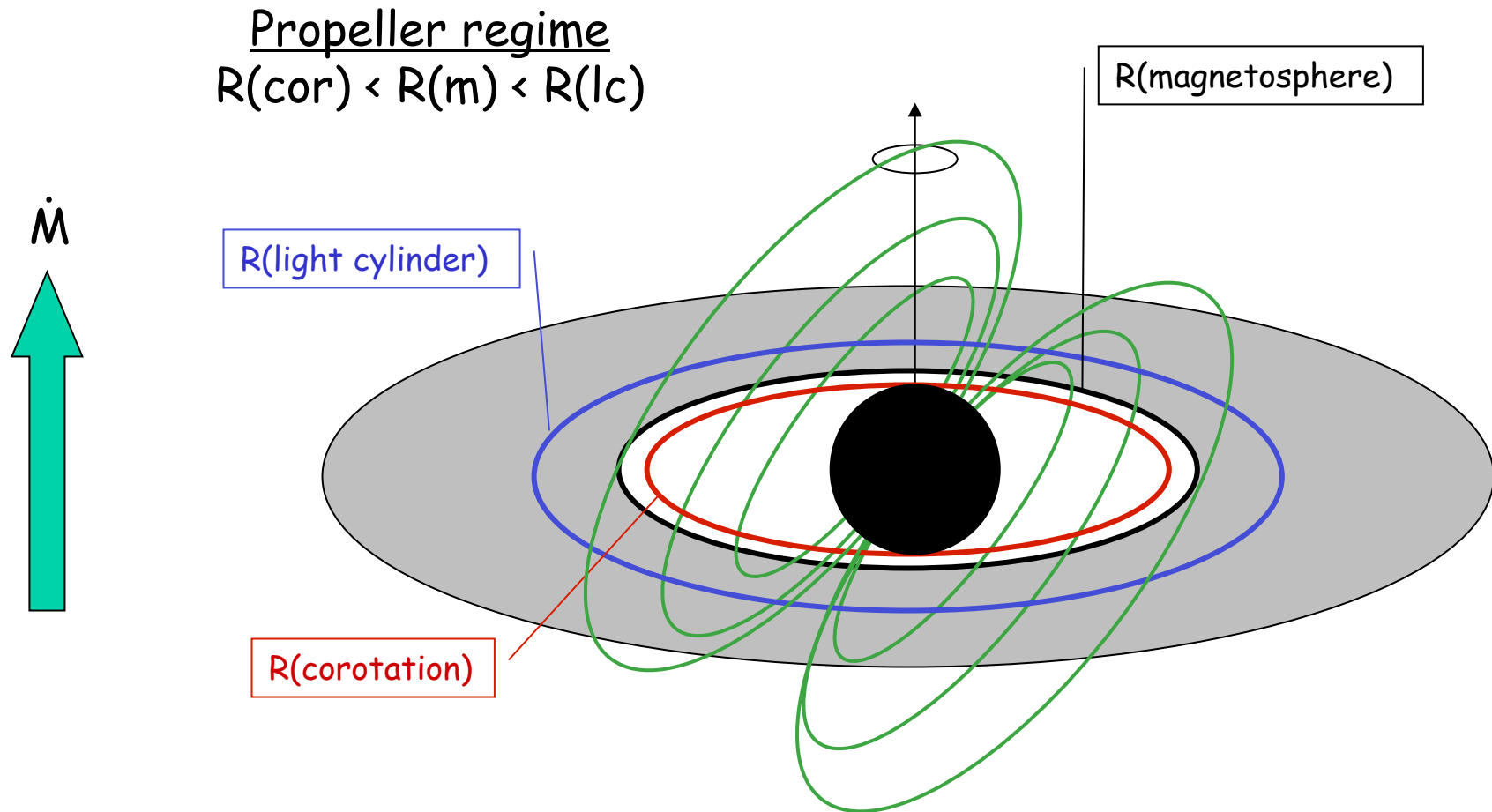
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Onset of centrifugal barrier in Hard XRTs

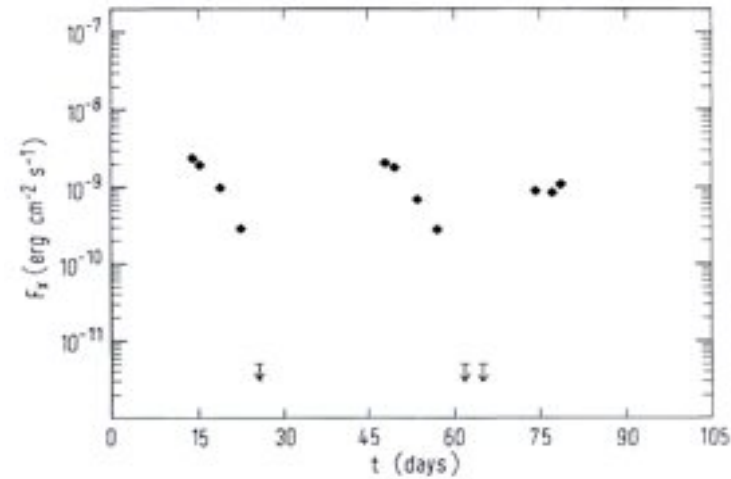
Very sharp X-ray luminosity decrease close to outburst end

V0332+53

$P=4.4\text{ s}$, $B_{\text{cyc}} \sim 10^{12}\text{ G}$

$L(\text{min}) \sim 10^{35}\text{ erg/s}$

(Stella, White & Rosner 1986)

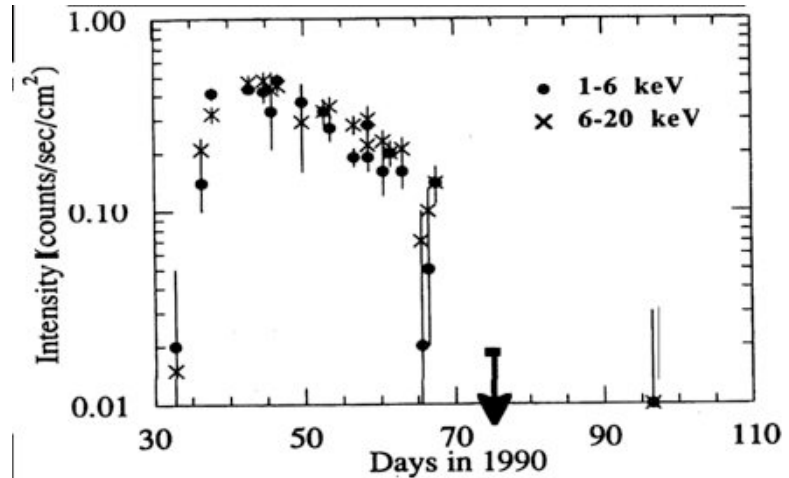


4U0115+63

$P=3.6\text{ s}$, $B_{\text{cyc}} \sim 10^{12}\text{ G}$

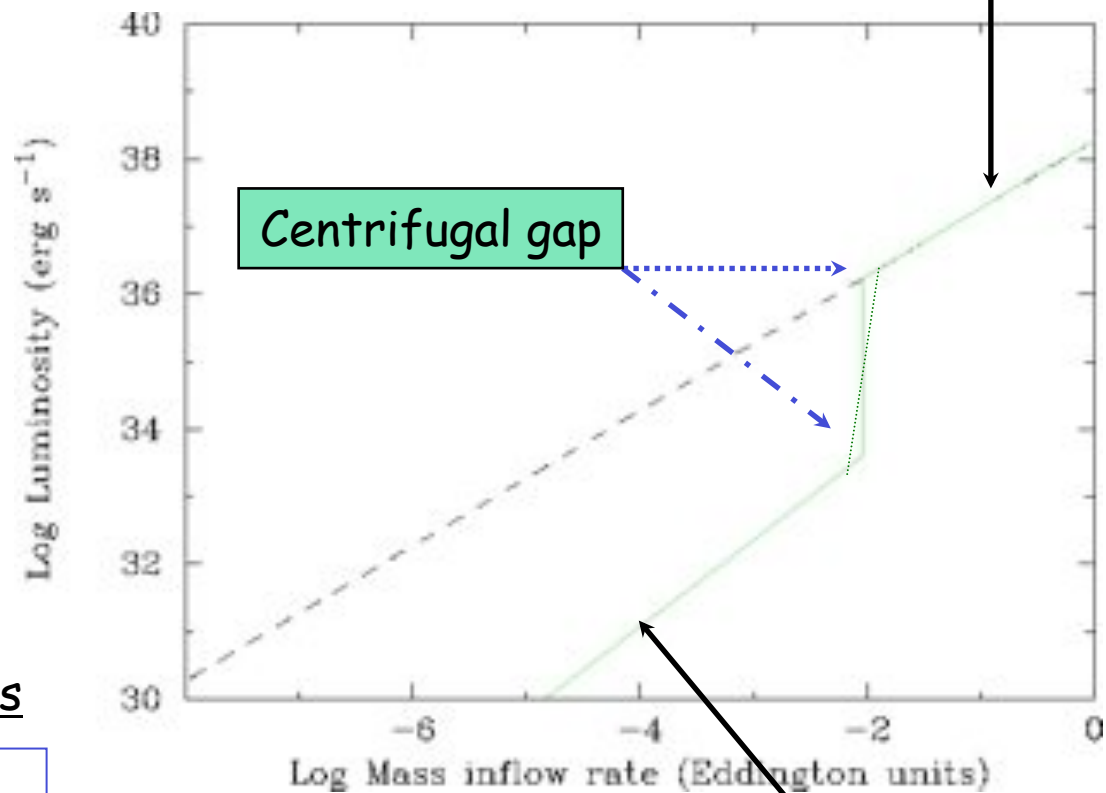
$L(\text{min}) \sim 10^{36}\text{ erg/s}$

(Tamura et al. 1992)



Expected mass-energy conversion efficiency in Hard XRTs

X-ray pulsar with $P \sim 4$ s $B \sim 10^{12}$ G



Centrifugal gap expectations

- factor of ~ 400 jump in L_x
- very steep dependence on $\dot{M}$ (not step-like)

(Stella et al. 1994; Corbet 1995)

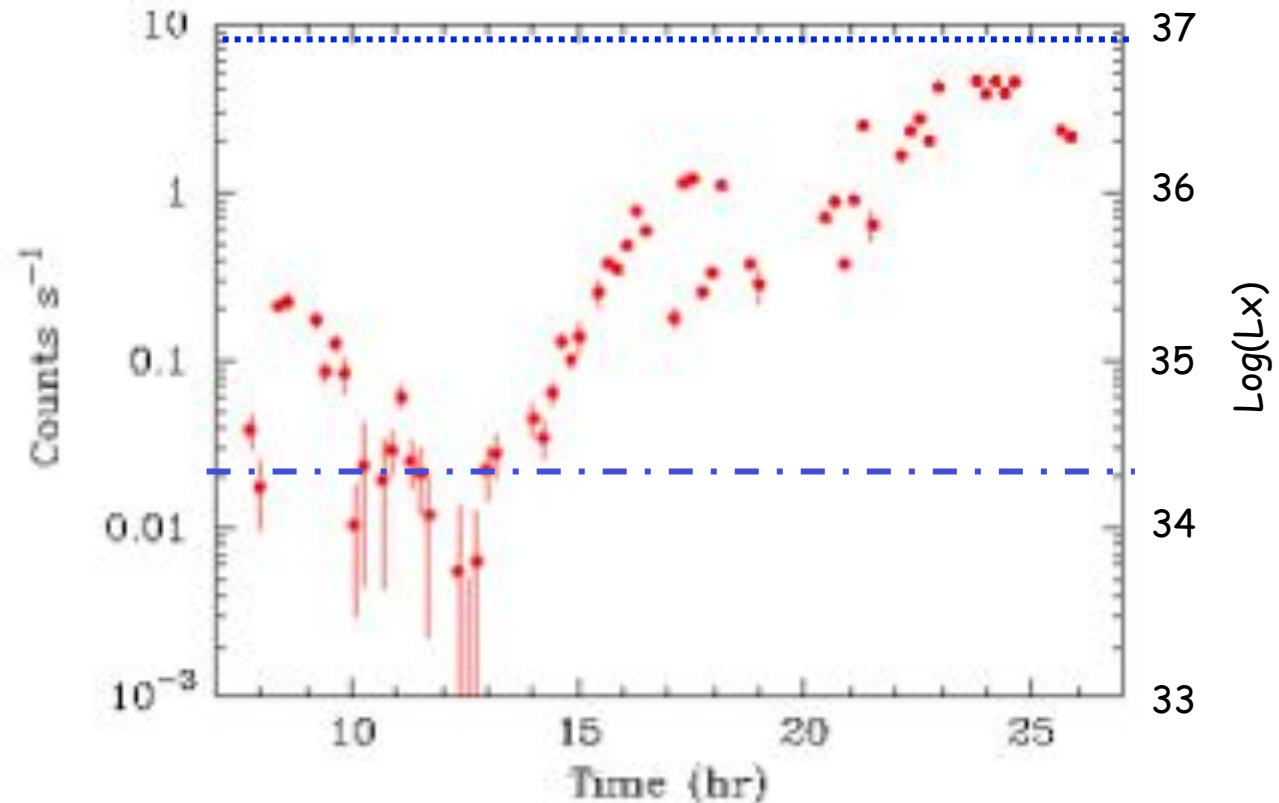
BeppoSAX observation 4U0115+63 in quiescence

(Campana et al. 2000)

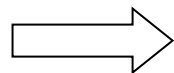


Expected range of centrifugal gap,
based on measured P, B and distance

- Observed L_x within centrifugal gap !
- Very Large (factor of ~ 250) L_x variations!
- 3.6 s pulsations present
- No substantial pulsation amplitude and spectral variations



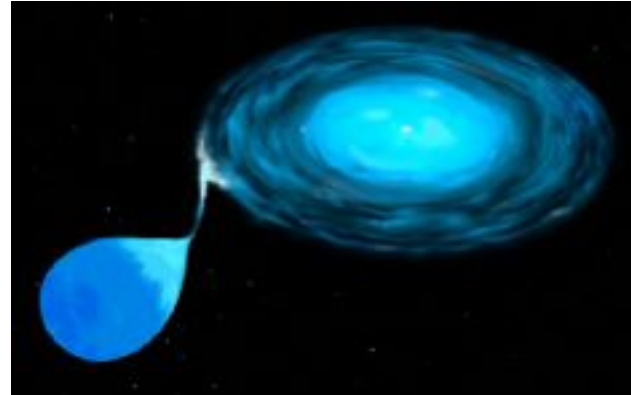
- Factor of < 3 variations in $\dot{M}$ expected at the most: imply a very steep dependence of L_x on $\dot{M}$ (power law slope of > 5)
- In the centrifugal gap some matter must leak through the barrier and accrete onto the NS surface



First evidence for centrifugal gap !

Soft XRTs: the missing link ?

Spin-up by accretion paradigm:
Soft XRTs (like NS low mass
X-ray binaries) contain old,
weakly magnetic NS with
millisecond periods



Soft XTRs should be among the progenitors
of ms radio pulsars

Quiescent Soft XRTs: $L_x \sim 10^{33}$ erg/s

Are there signs of a turned-on rotation powered
pulsar in quiescent Soft XRTs, when the mass inflow
rate is very low ?

➡ Nature of quiescent emission of Soft XTRs

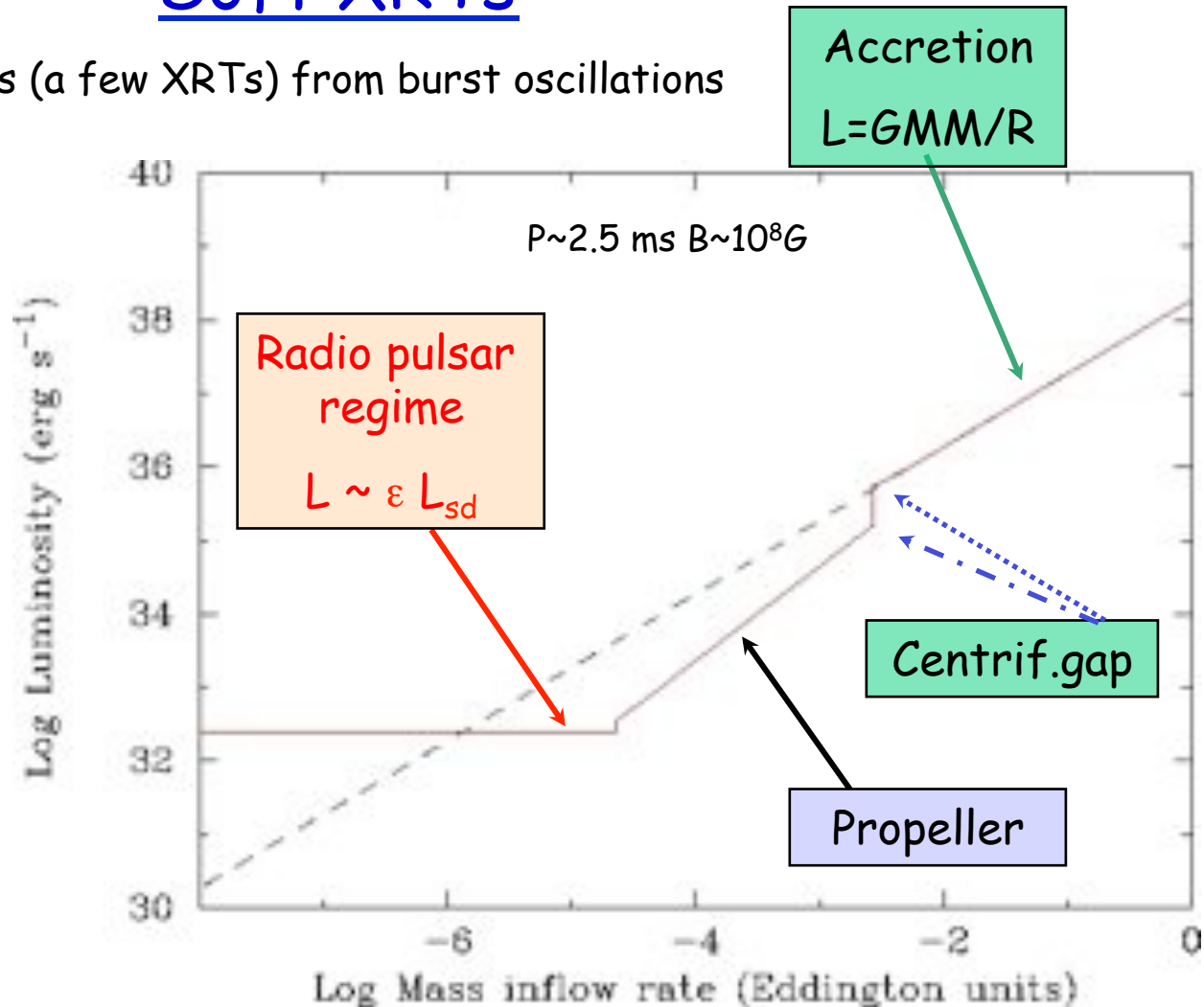
(Callanan 1989, Kulkarni et al. 1992, Stella et al. 1994)

Expected mass-energy conversion efficiency in Soft XRTs

$P \sim 1.6\text{--}4\text{ ms}$ in ~ 10 LMXRBs (a few XRTs) from burst oscillations

Basic expectations

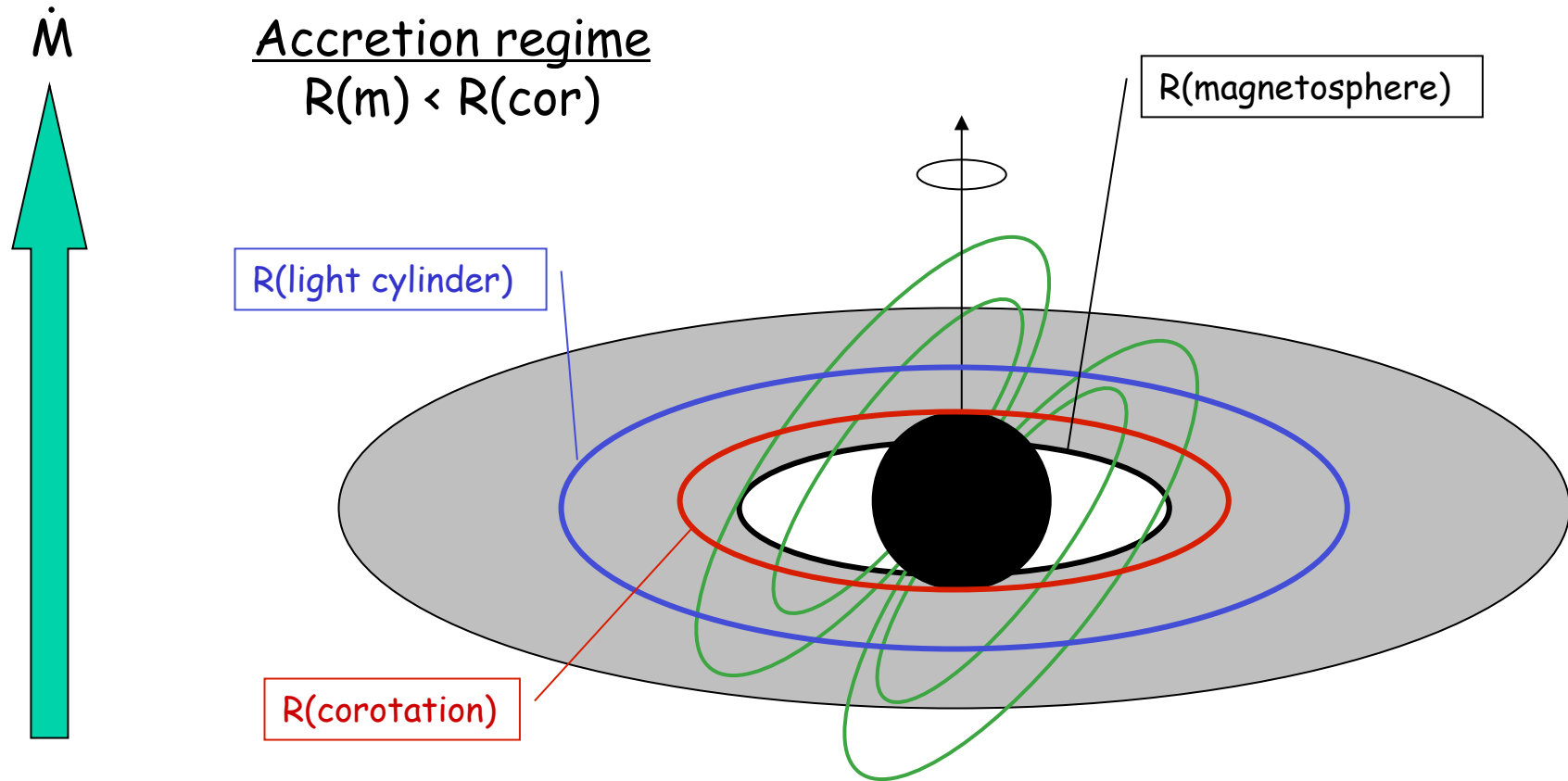
- small centrifugal gap
- propeller regime over a range of ~ 100 in $\dot{M}$
- radio pulsar regime for very low mass inflow rates: shock emission



(Stella et al. 1994,
Campana et al. 1998)

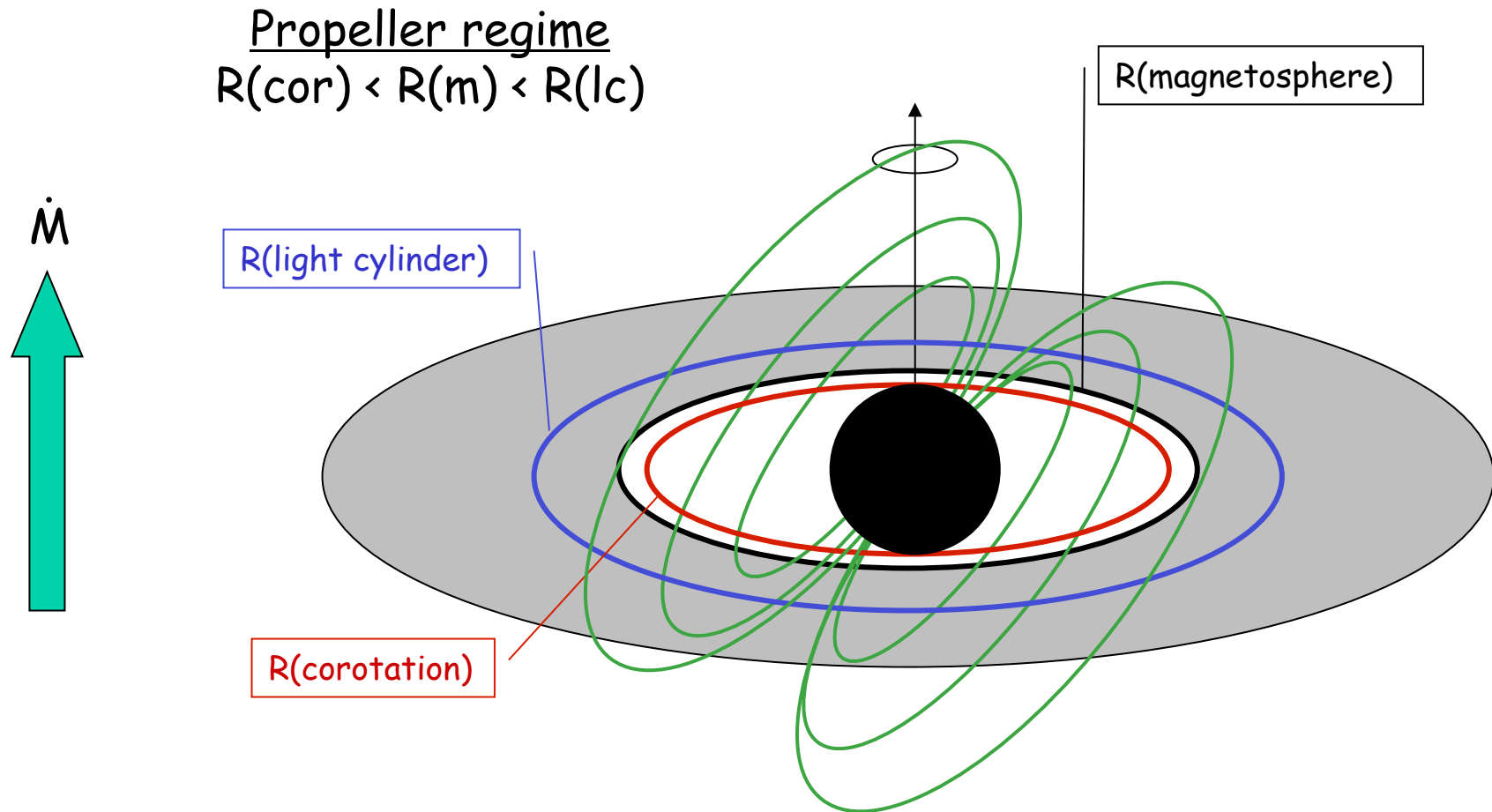
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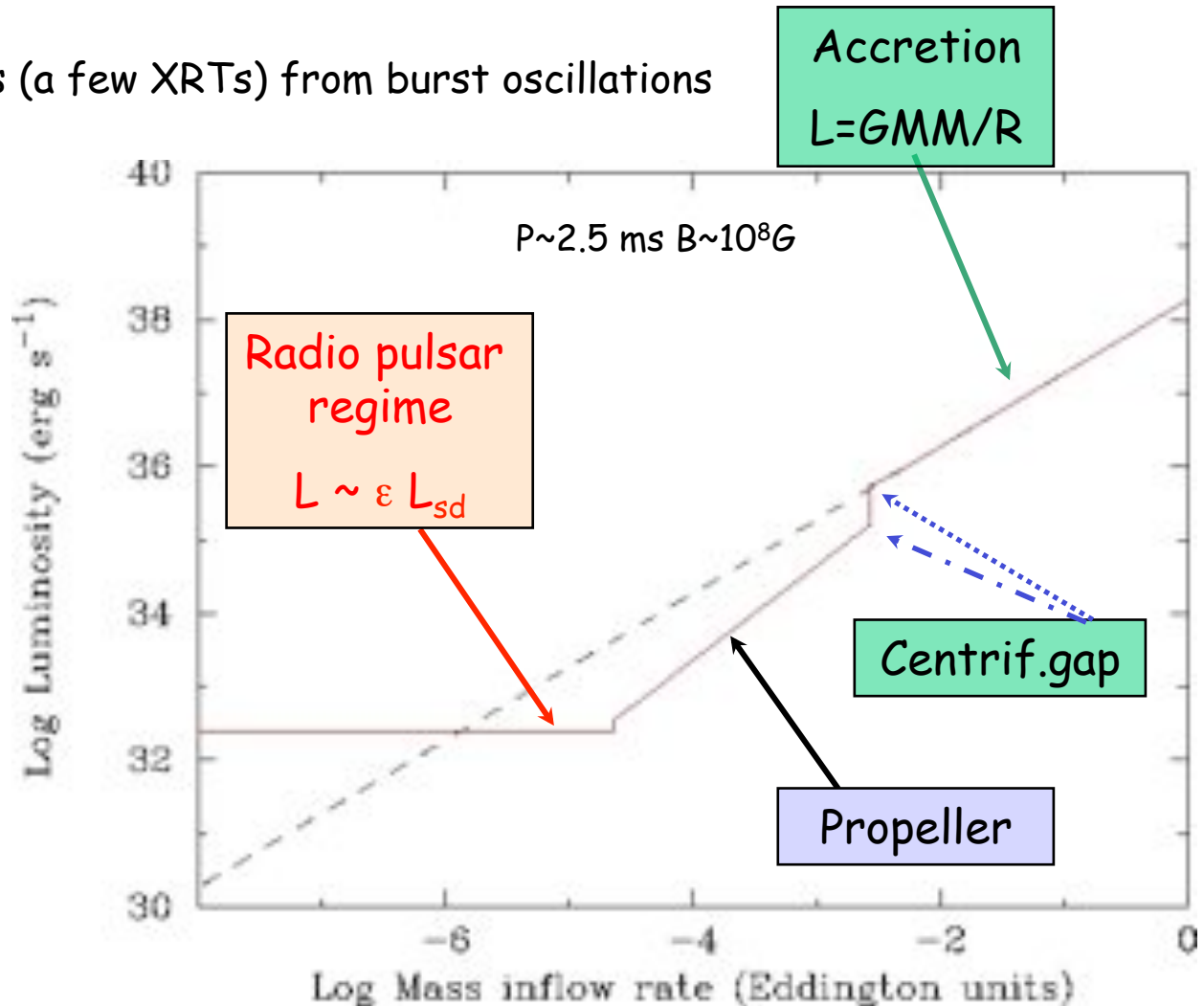
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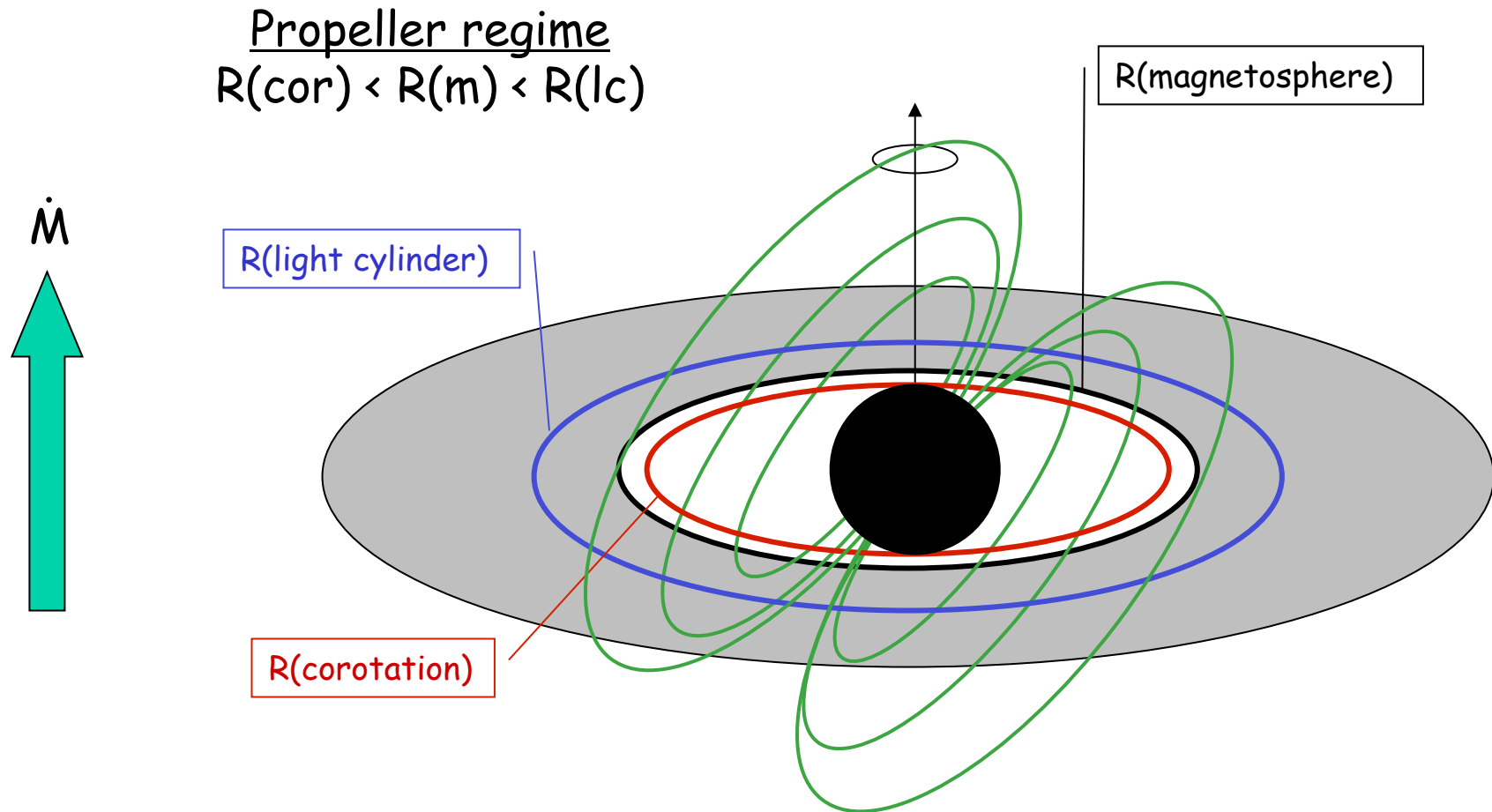
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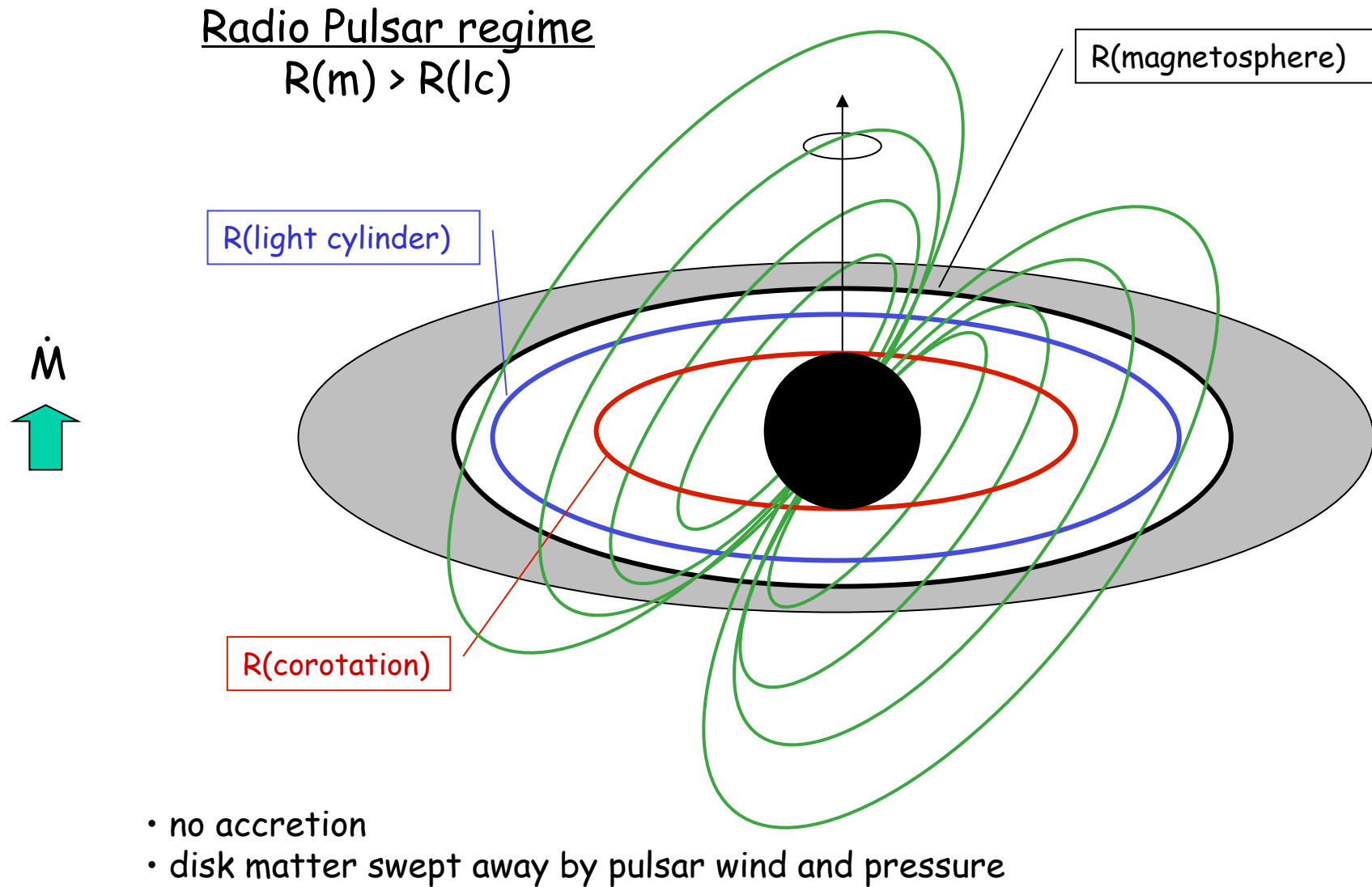
(Stella et al. 1994,
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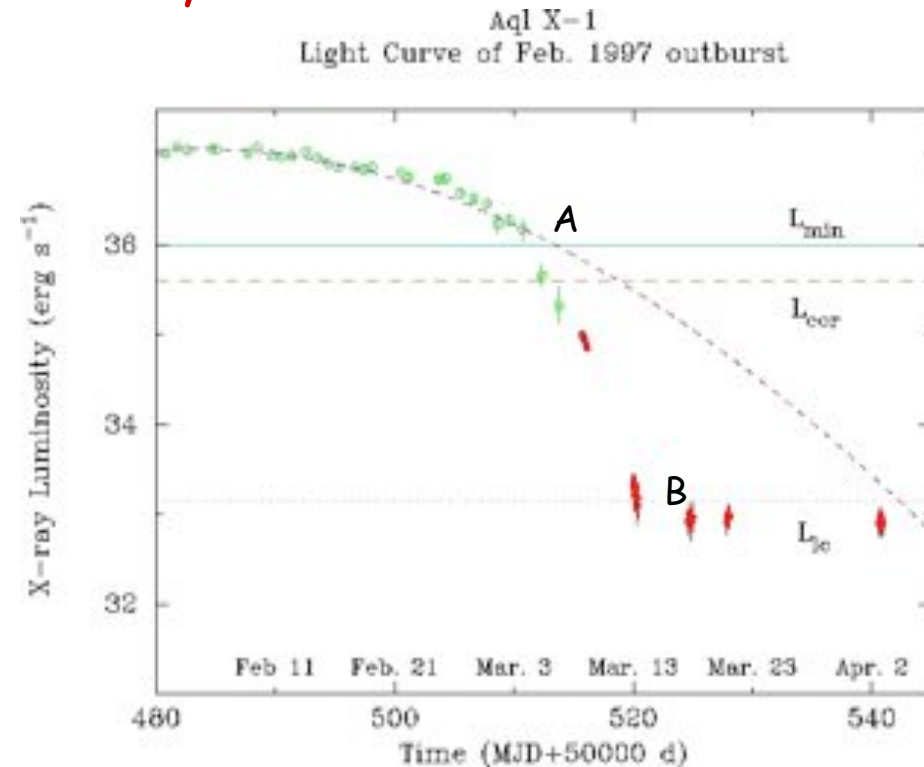


Soft XRTs: Aql X-1 from outburst to quiescence

Two transitions observed in outburst decay:

- A - $L_x \sim 10^{36}$ erg/s, decay ~ 1 d, spectrum hardens
- B - $L_x \sim 10^{33}$ erg/s, levels off, power law component decreases and flattens

$P = 1.8$ ms during bursts



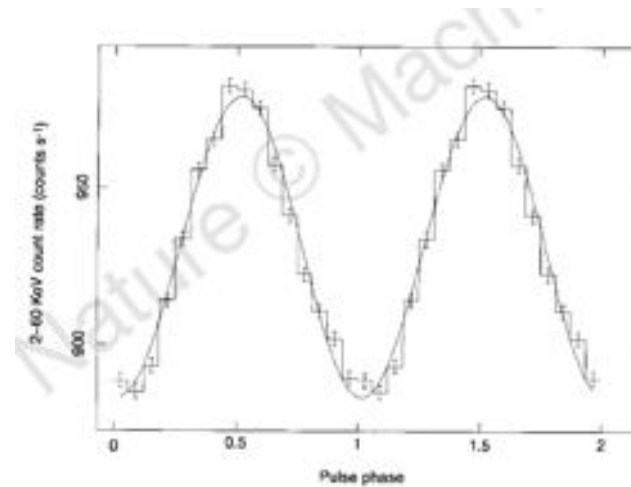
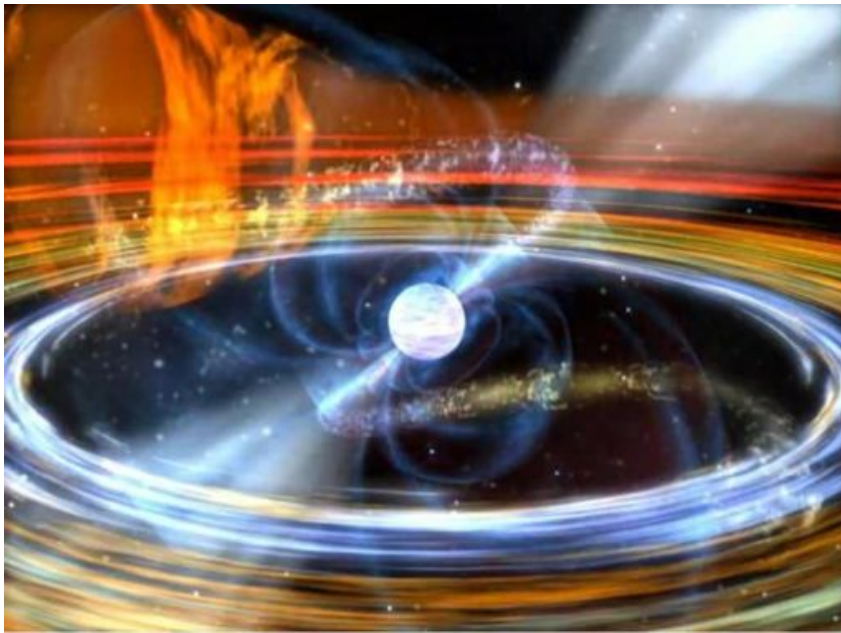
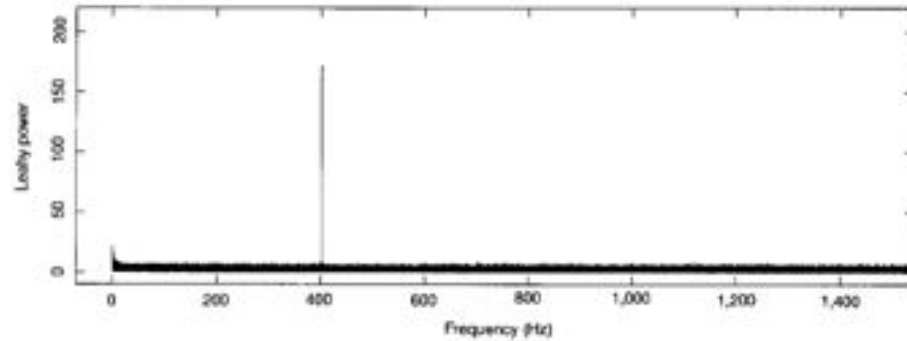
(Campana et al. 1998)

Discovery of accreting Millisecond pulsars

letters to nature

A millisecond pulsar in an X-ray binary system

Rudy Wijmands* & Michiel van der Klis*†



Discovery of accreting Millisecond pulsars

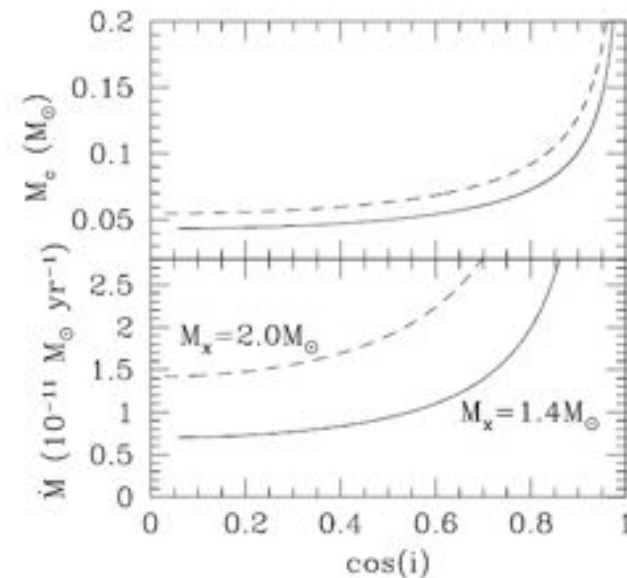
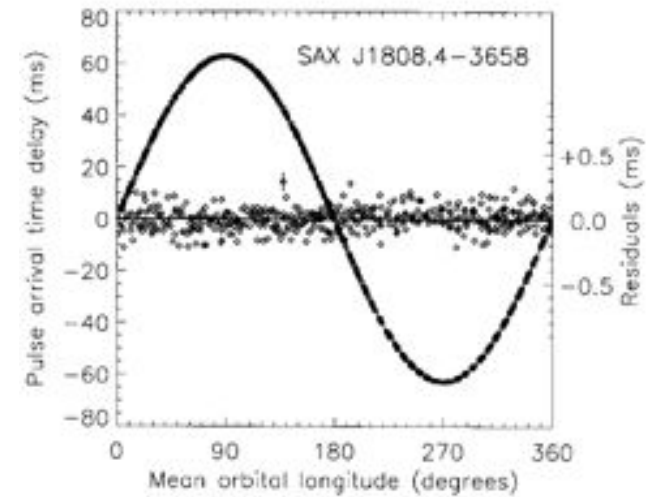
The two-hour orbit of a binary millisecond X-ray pulsar

Deepto Chakrabarty & Edward H. Morgan

Table 1 Observed parameters of SAX J1808.4-3658

Barycentric pulse frequency, ν_0 (Hz)	400.9752106(8)
Pulse frequency derivative, $ \dot{\nu} $ (Hz s^{-1})	$< 7 \times 10^{-13}$
Projected semimajor axis, $a \sin i$ (light ms)	62.809(1)
Orbital period, P_{orb} (s)	7249.119(1)
Epoch of 90° mean longitude, T_{90} (MJD)	50914.899440(1)
Eccentricity, e	$< 5 \times 10^{-4}$
Pulsar mass function, f_1 ($M_\odot$)	$3.7789(2) \times 10^{-4}$

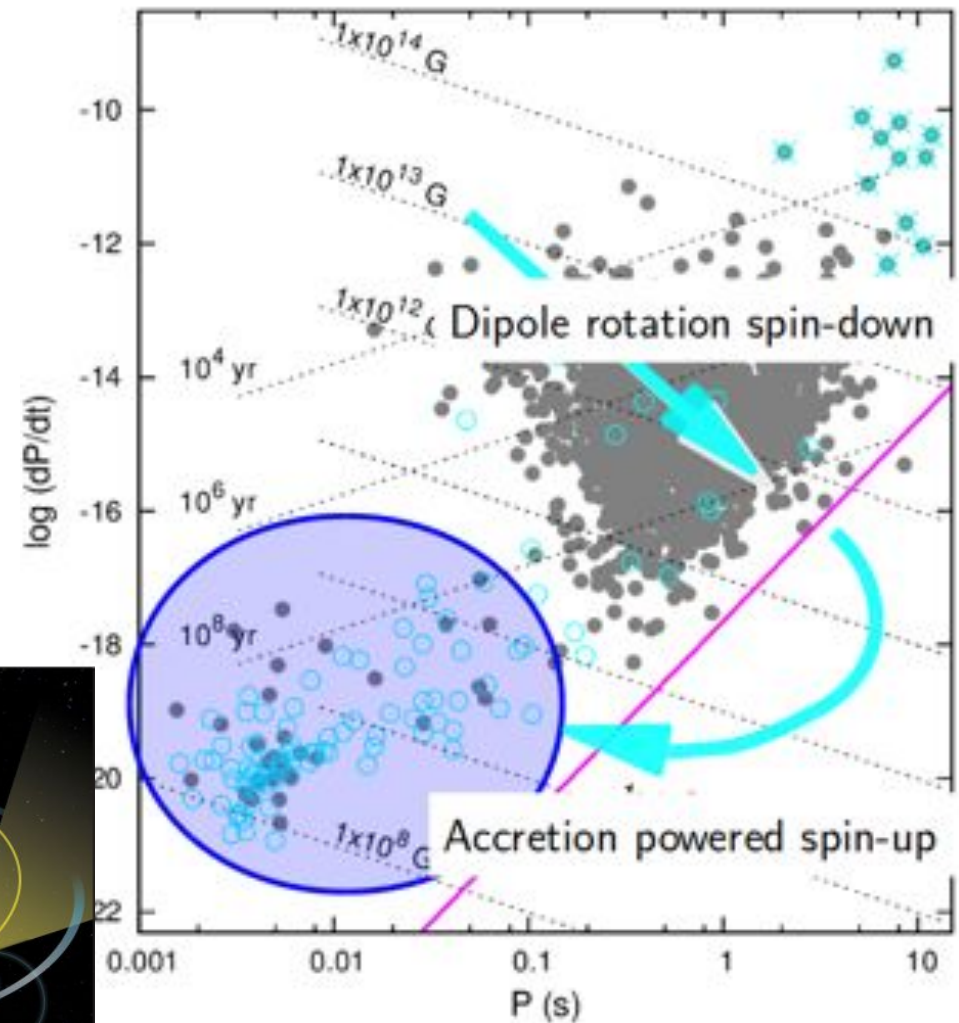
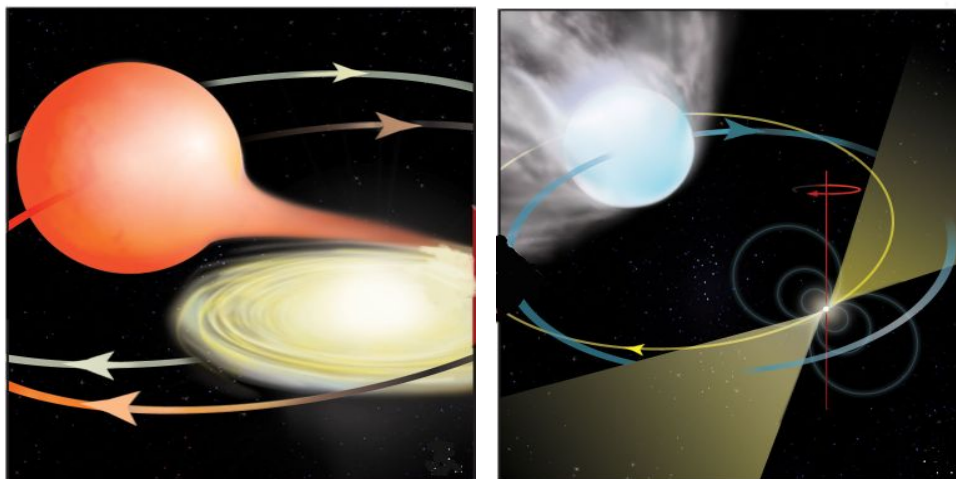
Modified Julian date MJD = JD - 2400000.5. Numbers in parentheses are the 1σ uncertainties in the last significant figure, and upper limits are quoted at the 2σ level. The ν limit refers to the magnitude, independent of sign. Epochs are given in Barycentric Dynamical Time (TDB).



The fundamental plane of pulsars

Accreting ms pulsars show that NS in low mass X-ray binaries are spun up to ms periods.

Do Accreting ms pulsars turn on as radio pulsars during quiescence?



A radio pulsar turning on in quiescence?

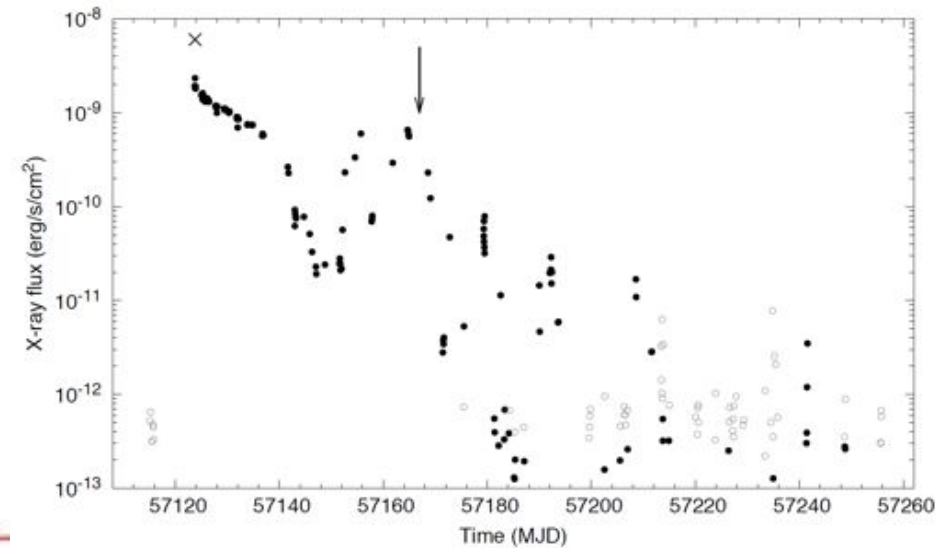
$$P_s = 1.5 \text{ ms}$$

$$R_{\text{co}} = [GM(P_s/2\pi)^2]^{1/3}$$

$$= 22 \text{ km}$$

$$\text{For } B \sim 10^8 \text{ G} \rightarrow \mu \sim 10^{26} \text{ G cm}^3,$$

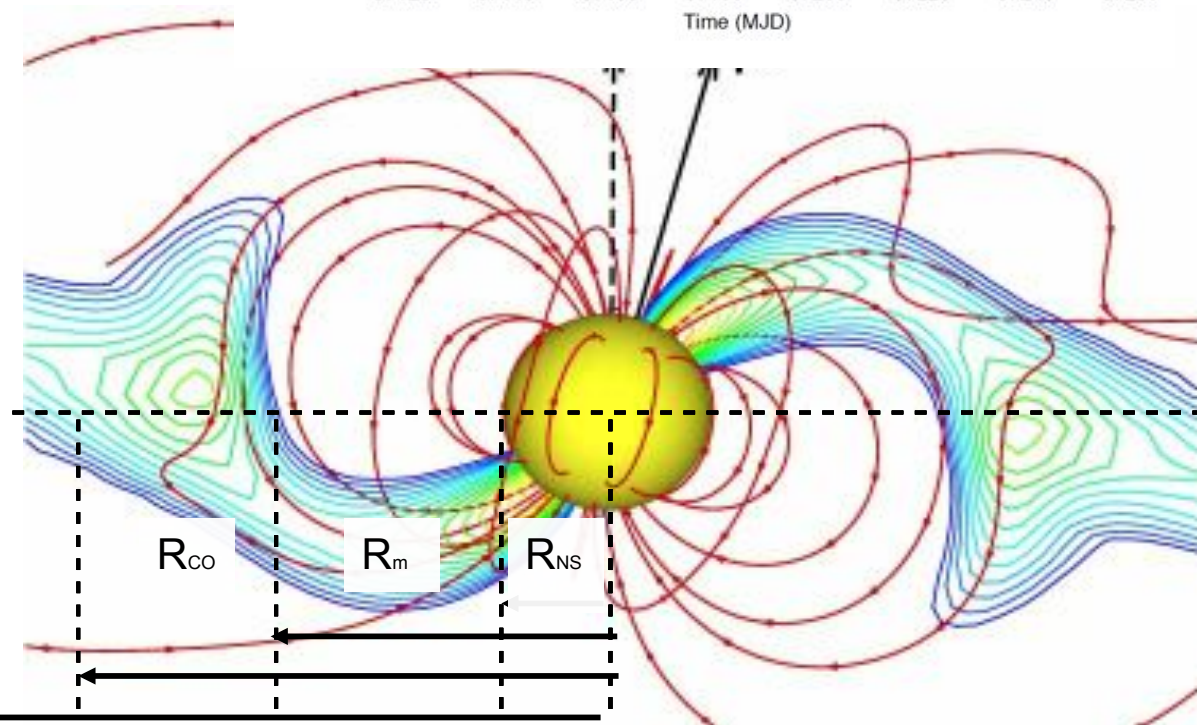
$$L_x = 10^{37} \text{ erg/s} \rightarrow \dot{M} = 5 \times 10^{16} \text{ g/s}$$



$$R_m = \xi \mu^{4/7} / M^{2/7} \\ (GM)^{1/7} = 20 \text{ km}$$

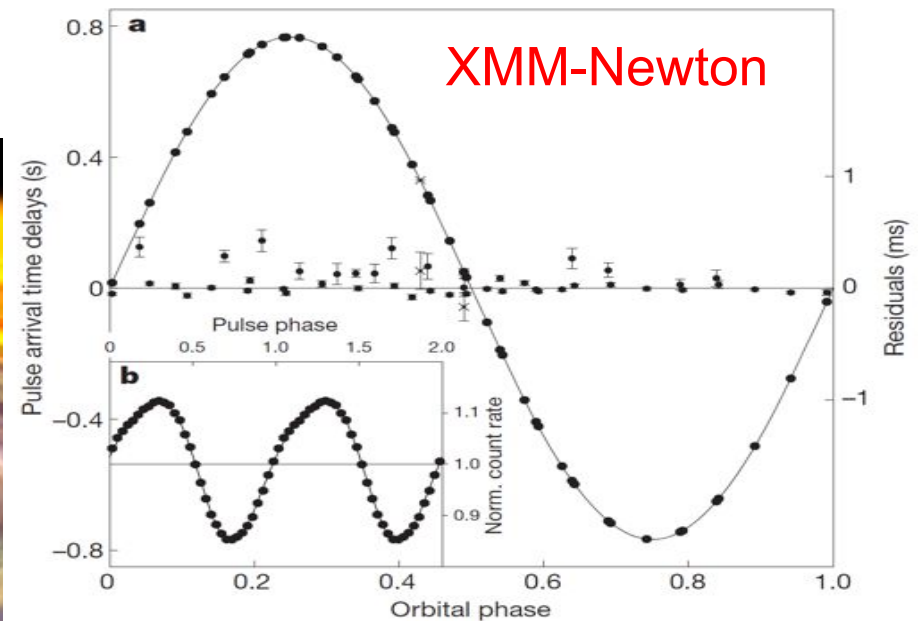
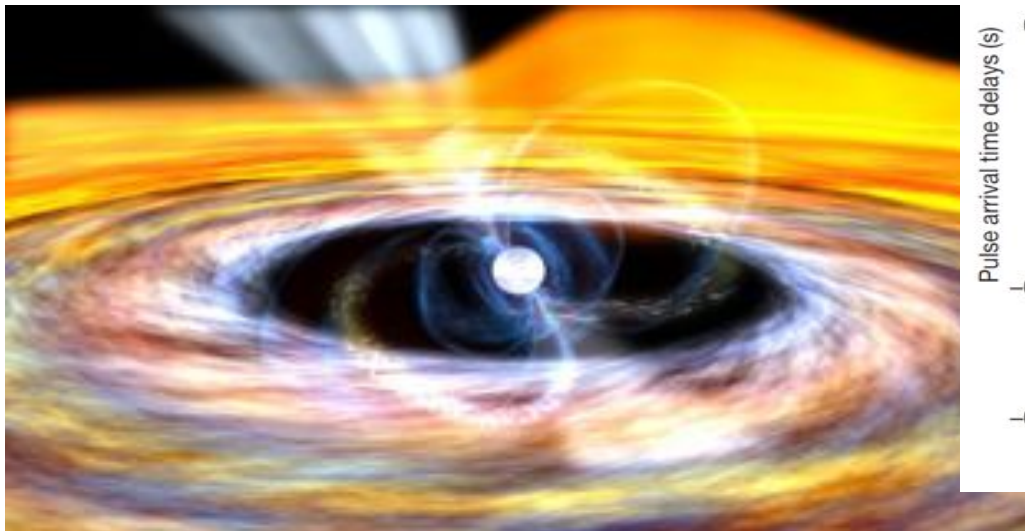
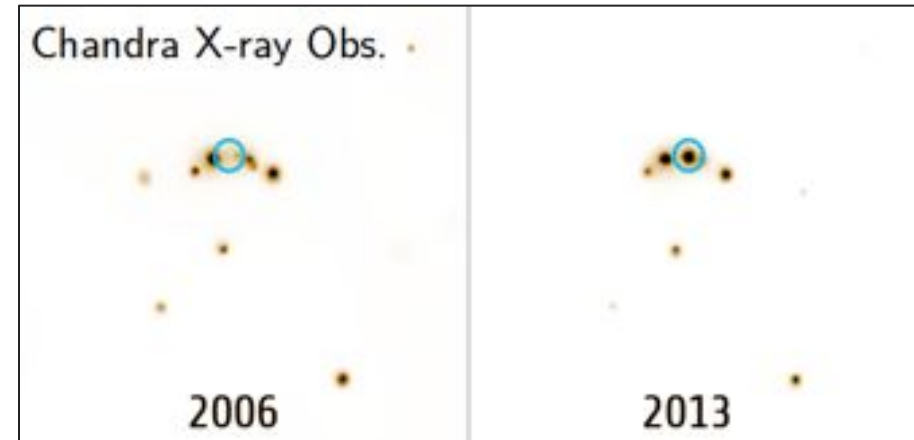
$$R_{\text{lc}} = cP_s / 2\pi = 72 \text{ km}$$

R_{lc}



An accreting ms pulsar in the globular cluster M28

IGR J18245-2452
 $L_{\text{(X-rays)}} \sim 5 \times 10^{36} \text{ erg/s}$
 $P_{\text{spin}} = 3.9 \text{ ms}$
 $P_{\text{orb}} = 11 \text{ hr}$

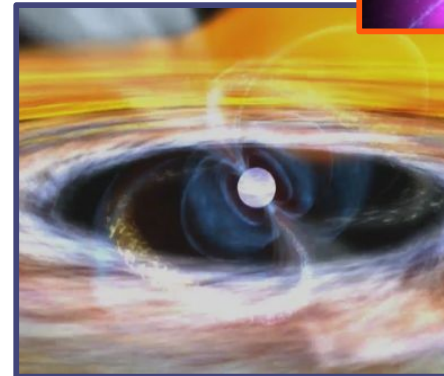
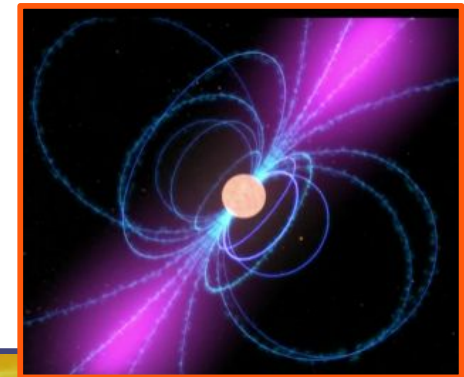


Swings between rotation and accretion power in a binary millisecond pulsar

Parameter	IGR J18245–2452	PSR J1824–2452I
Right Ascension (J2000)	18 ^h 24 ^m 32.53(4) ^s	
Declination (J2000)	−24° 52′ 08.6(6)″	
Reference epoch (MJD)	56386.0	
Spin period (ms)	3.931852641(2)	3.93185(1)
Spin period derivative	$< 2 \times 10^{-17}$	
RMS of pulse time delays (ms)	0.1	
Orbital period (hr)	11.025781(2)	11.0258(2)
Projected semi-major axis (lt-s)	0.76591(1)	0.7658(1)
Epoch of zero mean anomaly (MJD)	56395.216889(5)	

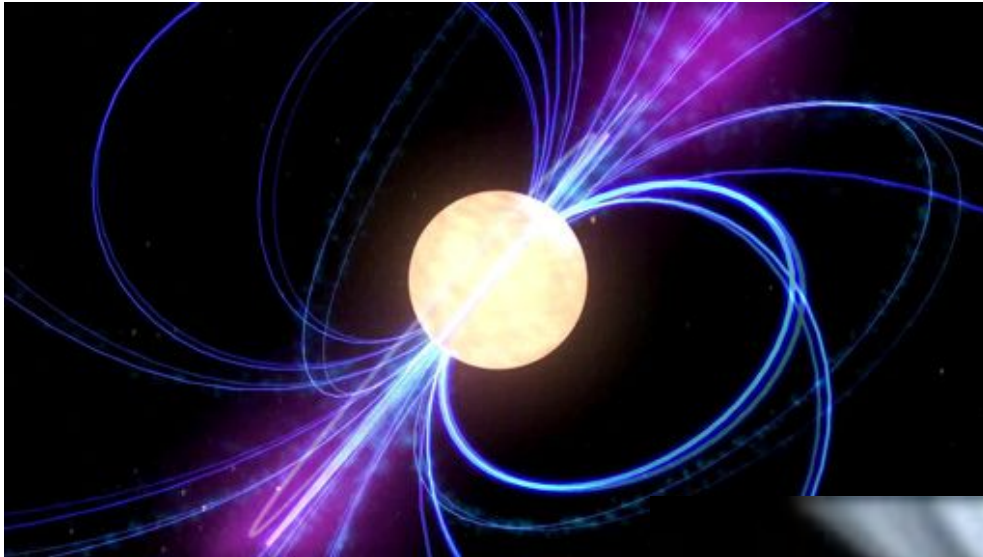
Papitto et al. 2013, Nature

Rotation powered **Radio PSR**



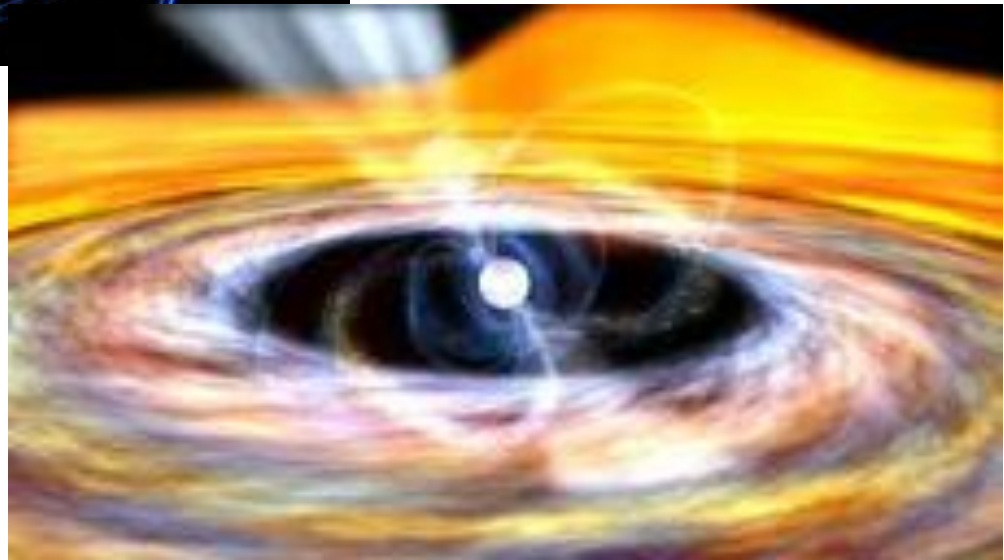
Xray PSR accretion powered

Transitional Millisecond Pulsars

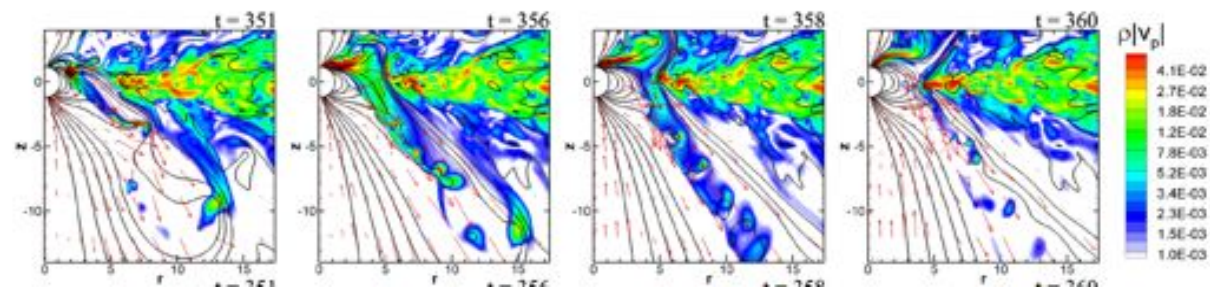
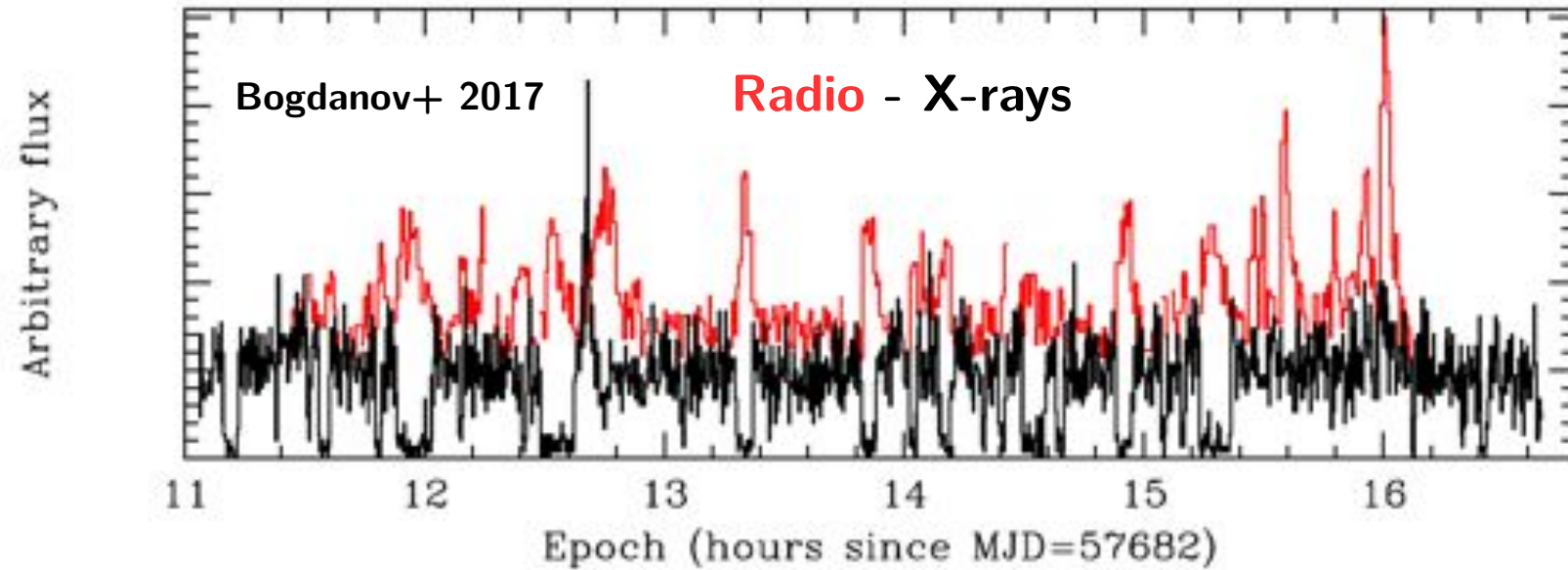


Rotationally-powered
(radio, gamma-ray)
pulsars

Accretion-powered
(X-ray) pulsars

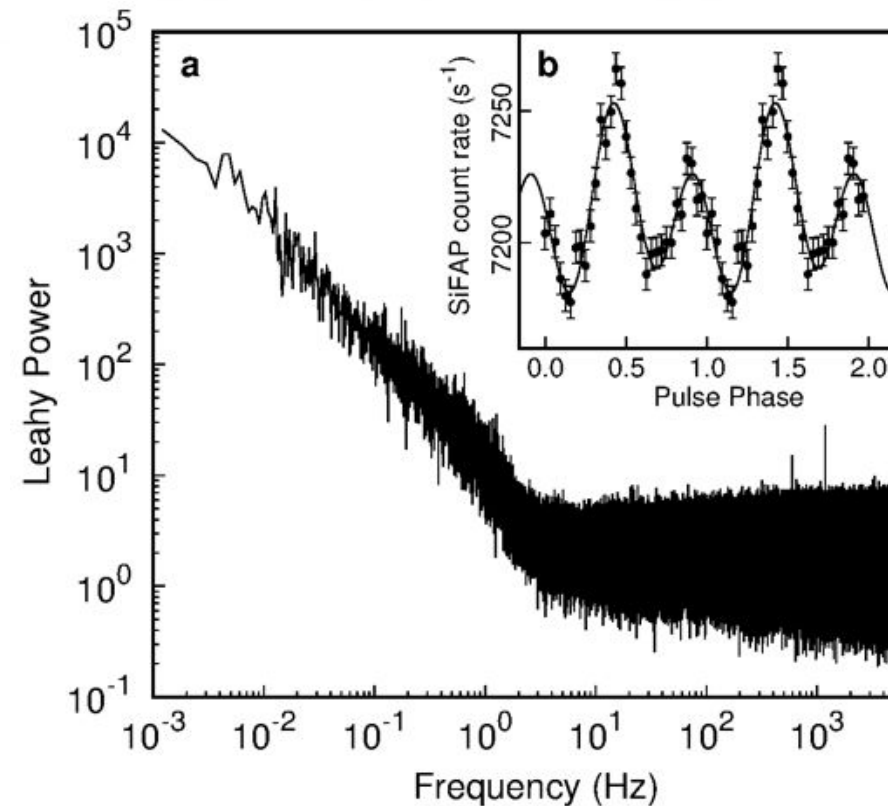


Plasma ejection events

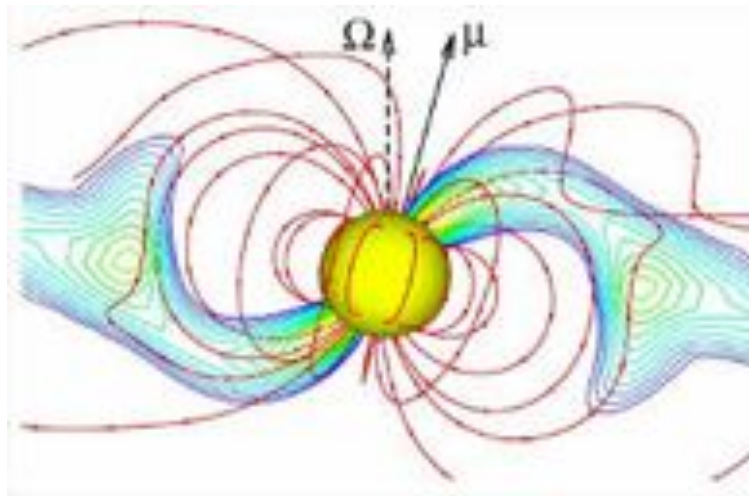


Optical pulsations from a transitional millisecond pulsar

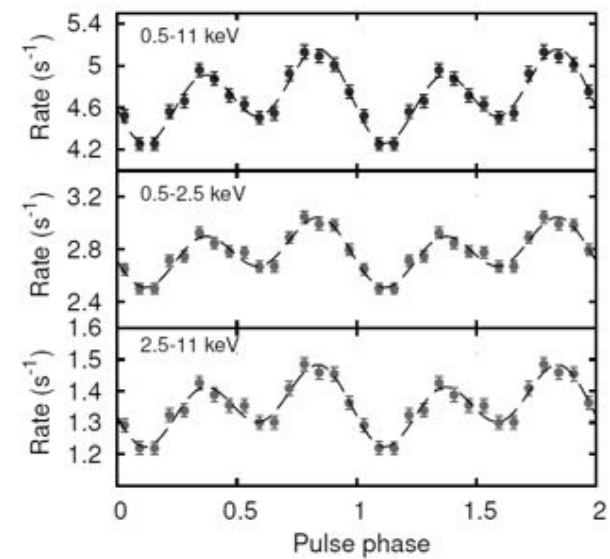
F. Ambrosino^{1,2}, A. Papitto^{3*}, L. Stella³, F. Meddi¹, P. Cretaro⁴, L. Burderi⁵, T. Di Salvo⁶, G. L. Israel³, A. Ghedina⁷, L. Di Fabrizio⁷ and L. Riverol⁷



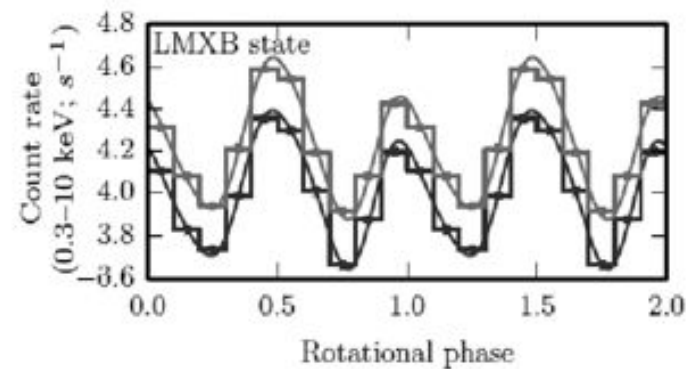
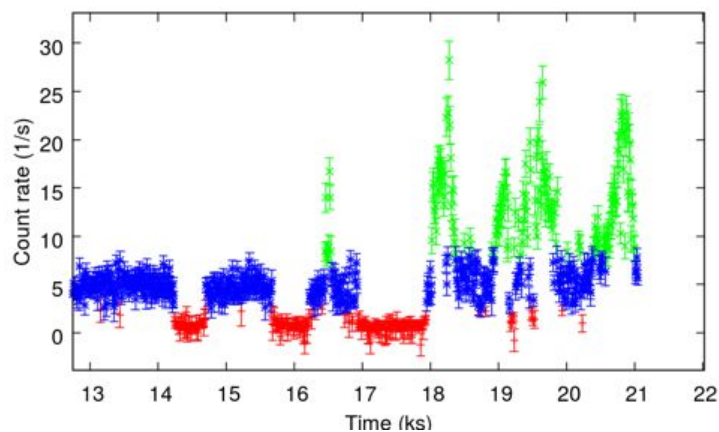
Coherent X-ray pulsations. Accretion power?



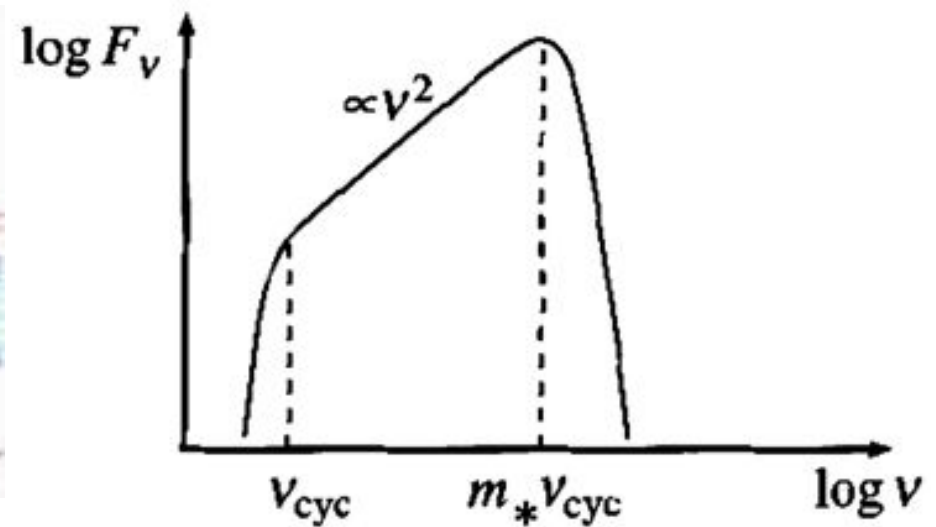
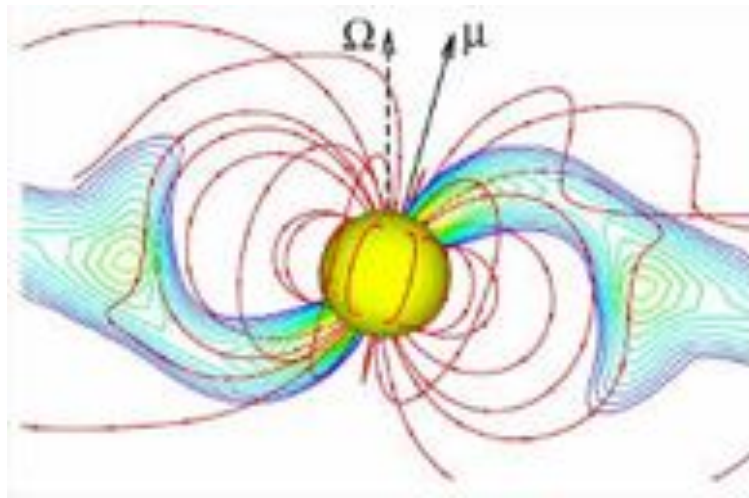
XSS J12270-4859 [Papitto+ 2015, MNRAS]



PSR J1023+0038 [Archibald+ 2015, ApJ]



Cyclotron emission from the hot-spots?

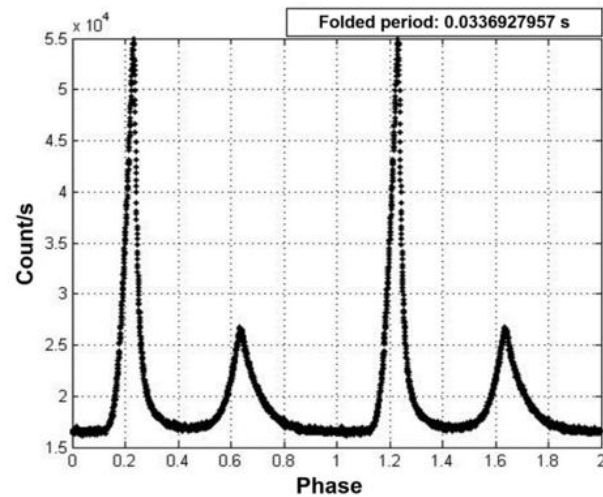


$$E_{\text{cyc}} \simeq 1 (B_s/10^8 \text{ G}) \text{ eV}$$

$$L_{\text{cyc}} \simeq 3 \times 10^{29} \left(\frac{A}{100 \text{ km}^2} \right) \left(\frac{kT_{\text{el}}}{100 \text{ keV}} \right) \text{ erg s}^{-1}$$

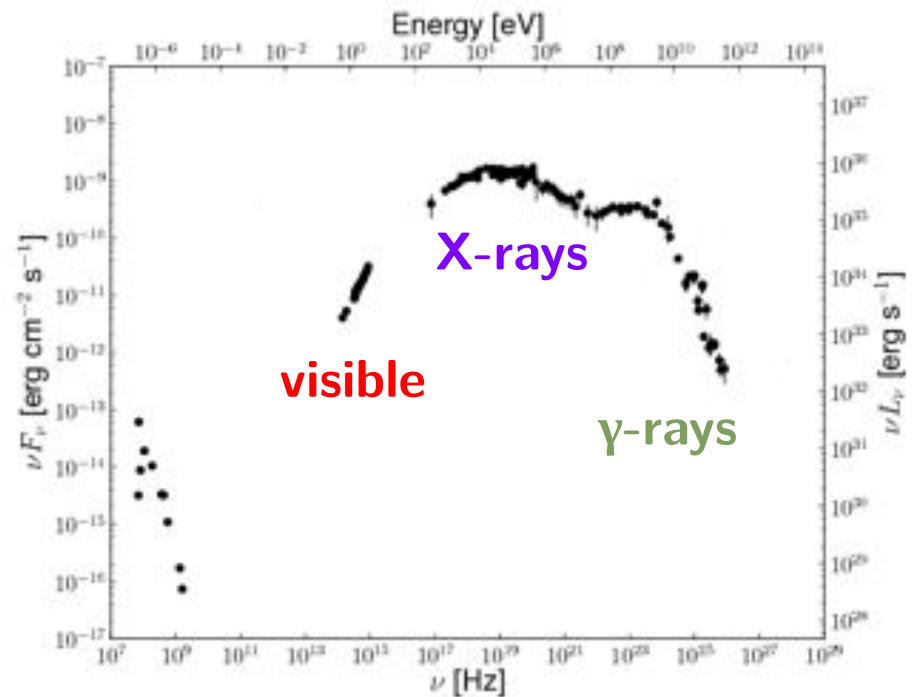
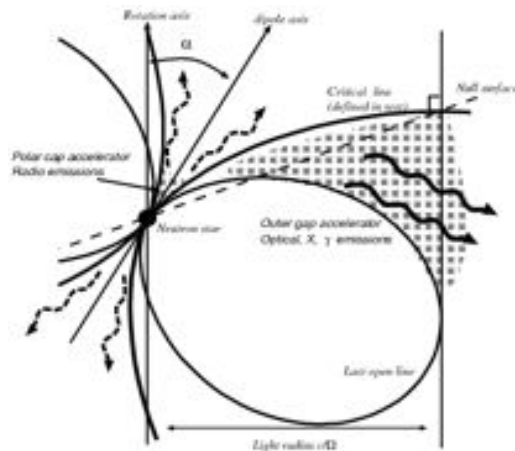
~30 times lower than observed luminosity even for very large spots

Rotationally-powered optical pulsations



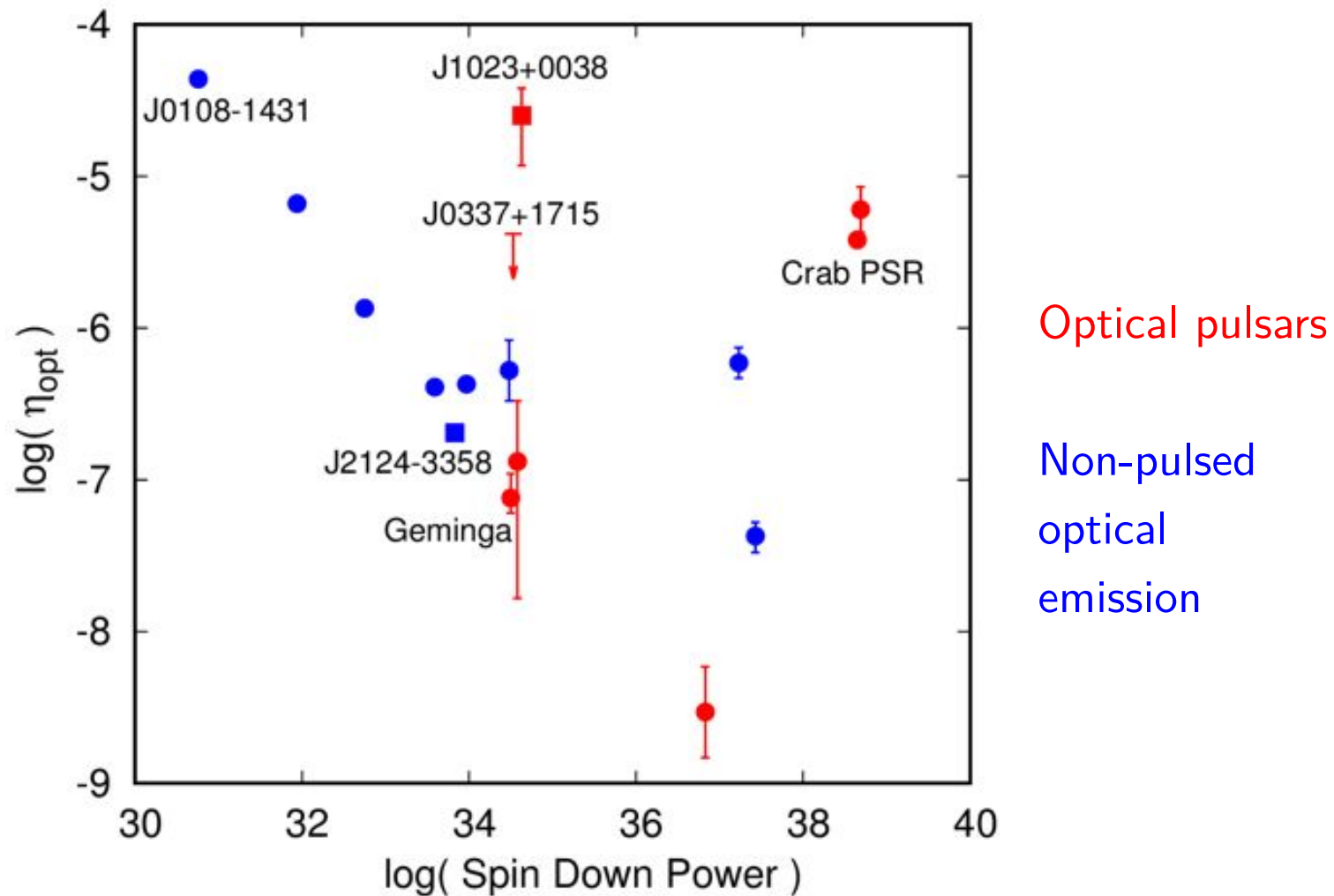
Crab/Vela/Geminga

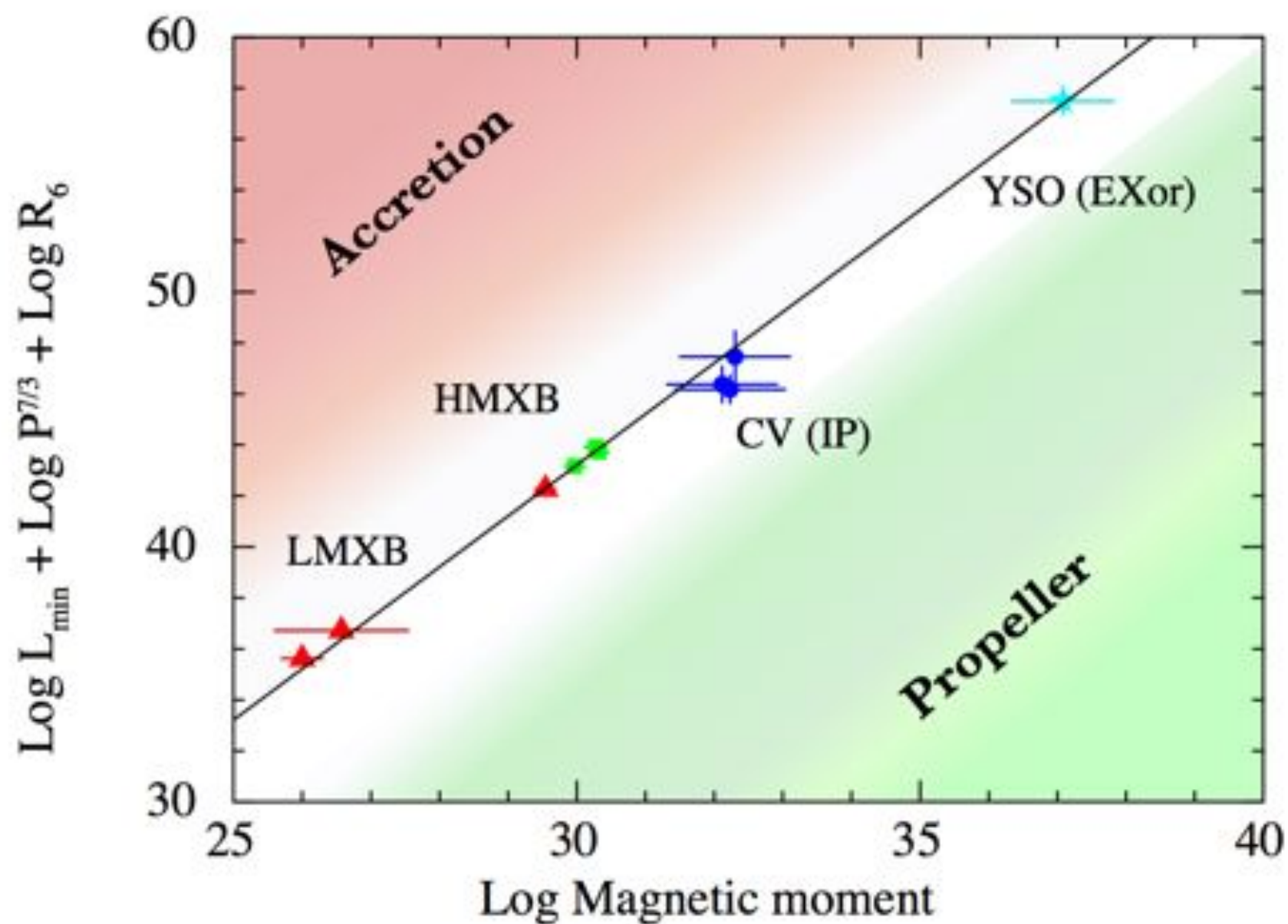
Synchrotron emission of $\sim$ MeV electrons



A very efficient optical rotationally-pwd pulsar

$$L(\text{B-band})/L_{\text{spin-down}} = 5 \times 10^{-4}$$





Propeller Onset

Simple theory

$$L_{\text{lim}} \sim 1.97 \times 10^{38} \xi^{7/2} \mu_{30}^2 P^{-7/3} M^{-2/3} R_6^{-1}$$

From observations

$$r_m = k (\mu^\alpha L^{-\gamma} R^{-\gamma})^{\frac{2}{3\beta}}, \quad (5)$$

which is independent of the central mass object. For standard accretion theory values $\alpha = 2$, $\beta = 7/3$, and $\gamma = 1$ we derive

$$r_m = (3.3 \pm 0.5) \times 10^8 (\mu_{30}^4 L_{36}^{-2} R_6^{-2} M)^{\frac{1}{7}} \text{ cm}, \quad (6)$$

with L_{36} the accretion luminosity in units of $10^{36} \text{ erg s}^{-1}$. Eq. (3) instead becomes

$$L_{\text{lim}} \sim 1.65_{-0.60}^{+0.94} \times 10^{37} \mu_{30}^2 P^{-7/3} M^{-2/3} R_6^{-1} \text{ erg s}^{-1}. \quad (7)$$

Need an advanced theory of disk–magnetosphere interaction

Two classes of high energy sources contain magnetars

(They are in highly absorbed regions in the galactic plane)

	<u>Soft Gamma Repeaters</u> SGRs (1987)	<u>Anomalous X-ray Pulsars</u> AXPs (1995)
• No. of sources	~10	~10
• Spin Period	2-8 s	2-12 s
• Period Derivative	$\sim 10^{-11}$ s/s	$\sim 10^{-11}$ - 5×10^{-13} s/s
• Isolated	yes	yes
• Recurrent Bursts	$\sim 10^{38}$ - 10^{41} erg/s	$\sim 10^{37}$ - 10^{38} erg/s
• Giant Flares	yes	no
• Persistent emission	$\sim 10^{35}$ erg/s	$\sim 10^{34}$ - 10^{35} erg/s
• Association to SNRs	?	in some cases
• Radio pulsations	?	in some cases

SGR Bursts

- Concentrated in time ("outbursts")
- Relatively simple profiles (faster rise than decay)
- Broad distribution of wait times (~ 7 decades) : similar to that of earthquakes; no waiting-time correlations
- Energy distribution $dN/dE \sim E^{-5/3}$: similar to earthquakes
- Most bursts release $\sim 10^{38}$ - 10^{41} ergs

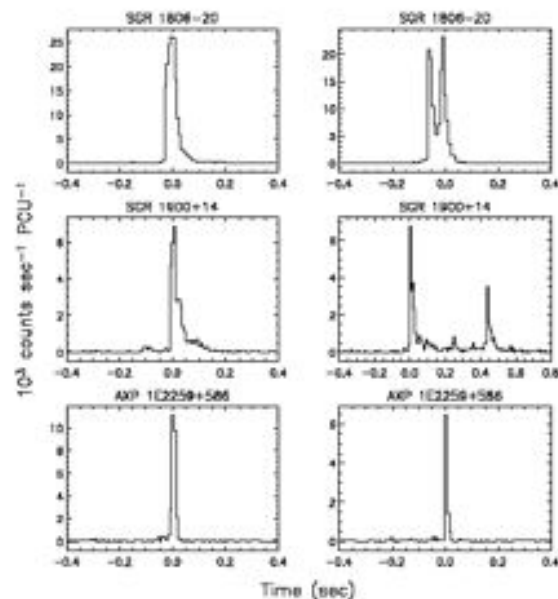
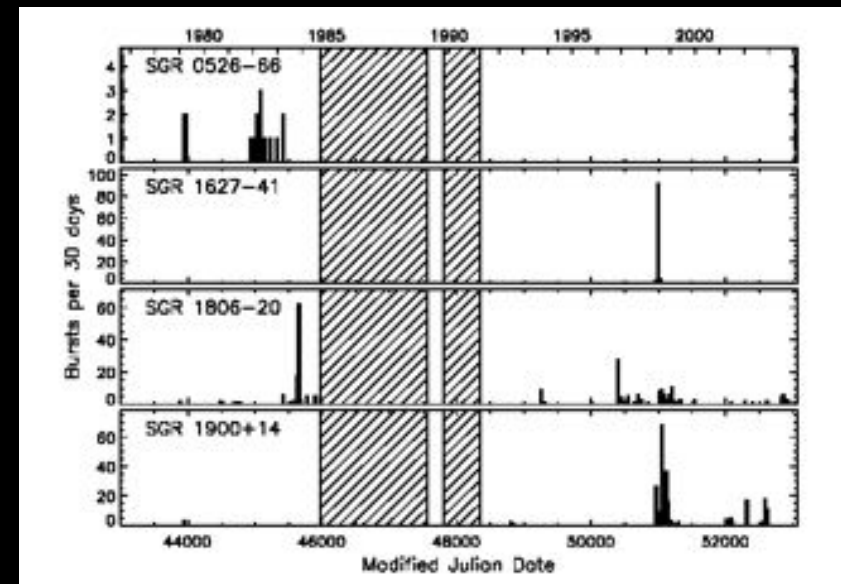
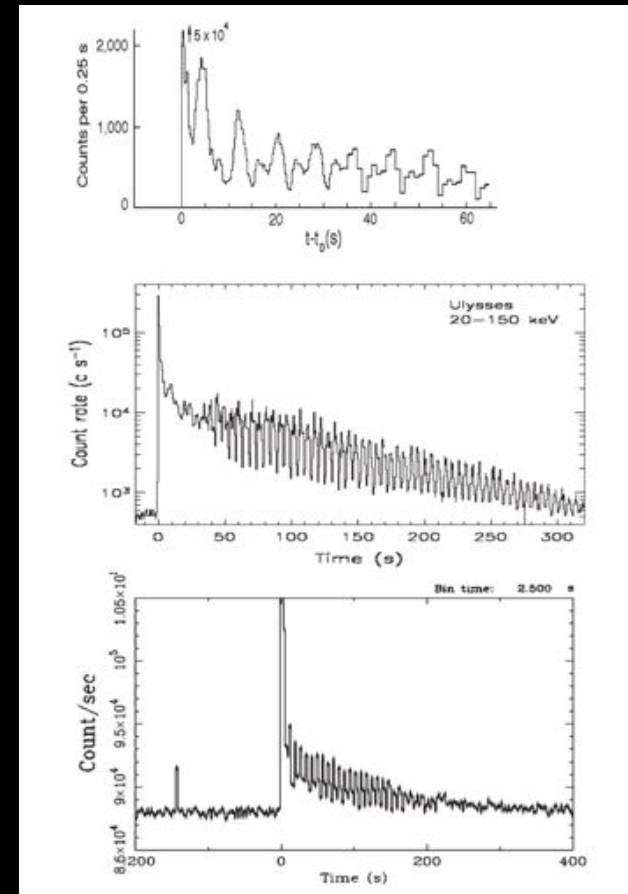
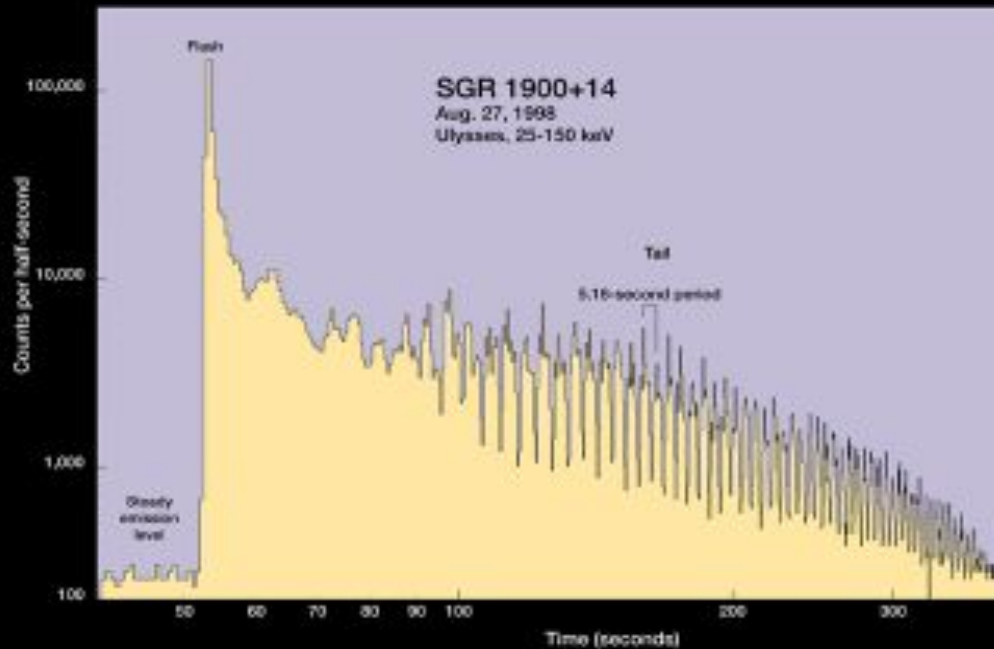


Fig. 14.1. A selection of common burst morphologies recorded from SGR 1806-20, SGR 1900+14 and 1E 2259+586, as observed with the *RXTE* PCA. All light curves display counts in the energy range 2-20 keV, with a time resolution of 7.8 ms. See text for further details.



Giant Flares

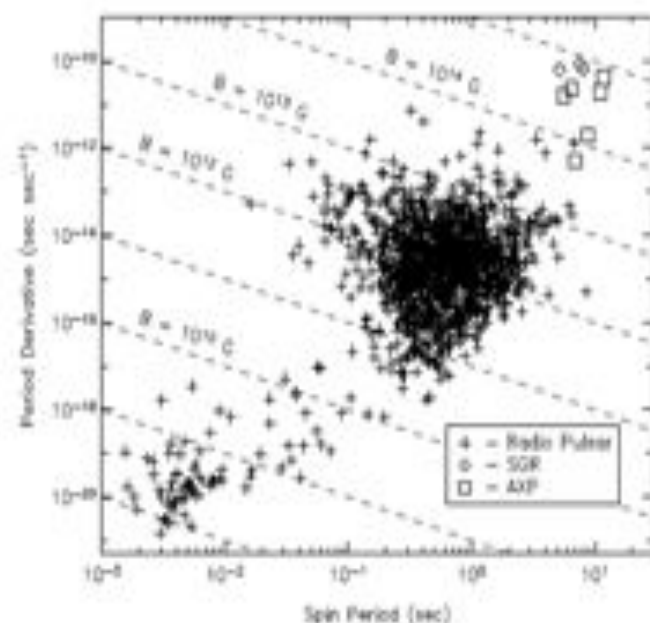
- 1979 March 5 from SGR 0526-66: energy released $\sim 10^{44}$ ergs
- 1998 August 27 from SGR1900+14: energy released $\sim 10^{44}$ ergs
- Main peak ~ 0.2 s; variability down to ms, hard BB-like spectrum ($kT \sim 100$ keV)
- 2-3 min long ringing tail: (nearly) exponential decay
pulsations at 5-8 s (spin period)
TB temperature of ~ 30 -50 keV.
- In Magnetar model: GFs due to large scale rearrangements of core B-field



- Period range: 5-12 s
- Characteristic ages: 10^3 - 10^5 yr

RAPID SPIN DOWN

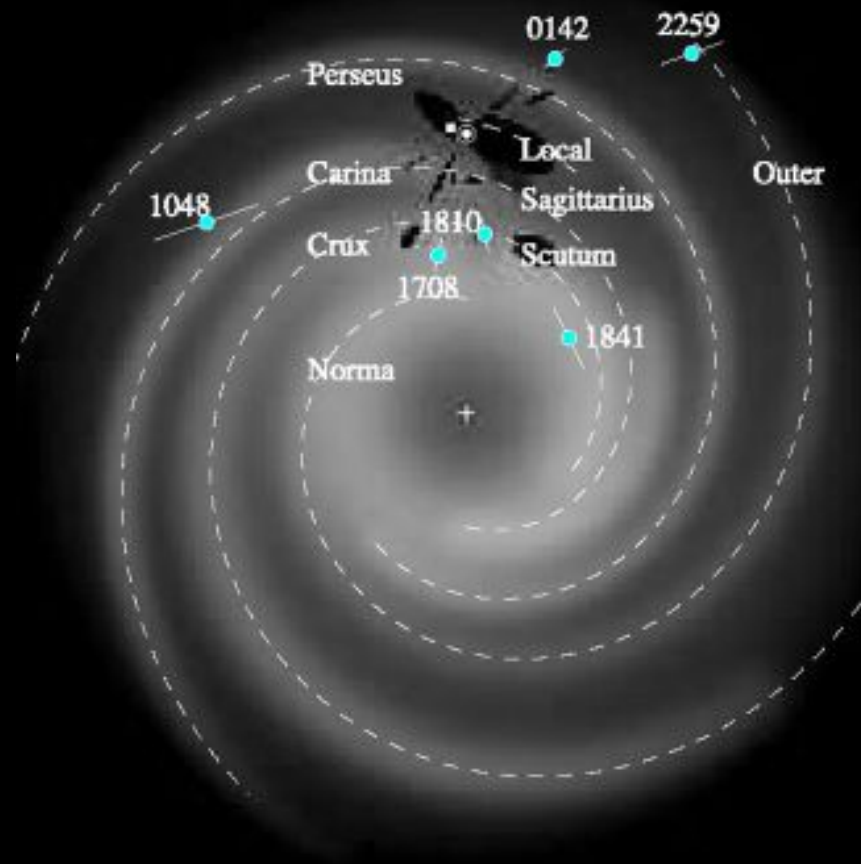
- Estimated $B \sim 10^{14}$ - 10^{15} G (see table)



Source	Period (s)	Period Derivative (10^{-11} s s $^{-1}$)	Magnetic Field ^a (10^{14} Gauss)	Spin down Age ^b (10^3 years)
SGR 0526-66	8.0	6.6	7.4	1.9
SGR 1627-41	6.47	—	—	—
SGR 1806-20	7.5	8.3-47	7.8	1.4
SGR 1900+14	5.2	6.1-20	5.7	1.3
CXOU 010043.1-721134	8.0	—	—	—
4U 0142+61	8.7	0.20	1.3	70
1E 1048.1-5937	6.4	1.3-10	3.9	4.3
1RXS J170849-400910	11.0	1.9	4.7	9.0
XTE J1810-197	5.5	1.5	2.9	5.7
1E 1841-045	11.8	4.2	7.1	4.5
AX J1844-0258	7.0	—	—	—
1E 2259+586	7.0	0.048	0.60	220

- $B_{\text{dipole}} = 3.2 \times 10^{15} \sqrt{P\dot{P}}$ G (the mean surface dipole field)
- Characteristic age of pulsar spinning down via magnetic braking ($P/2\dot{P}$)

Magnetars in our Galaxy



Magnetars: Neutron Stars powered by magnetic energy

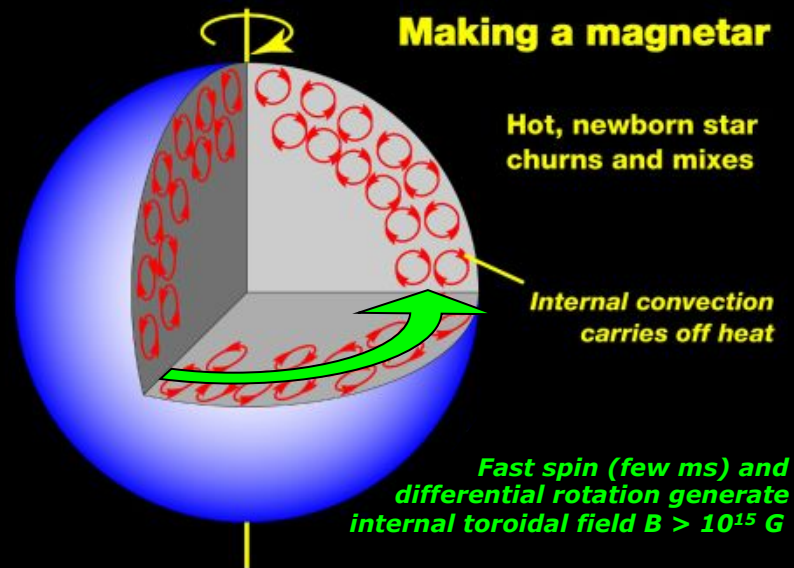
- "Magnetars" (MAGNETic sTARS): neutron stars with very high magnetic fields ($B > 10^{14} \text{G}$)
Why magnetic energy? (Magnetic energy propagating through fractures in the crust)

Persistent luminosity 10-100 times higher than spin down power

-> rotational energy ruled out

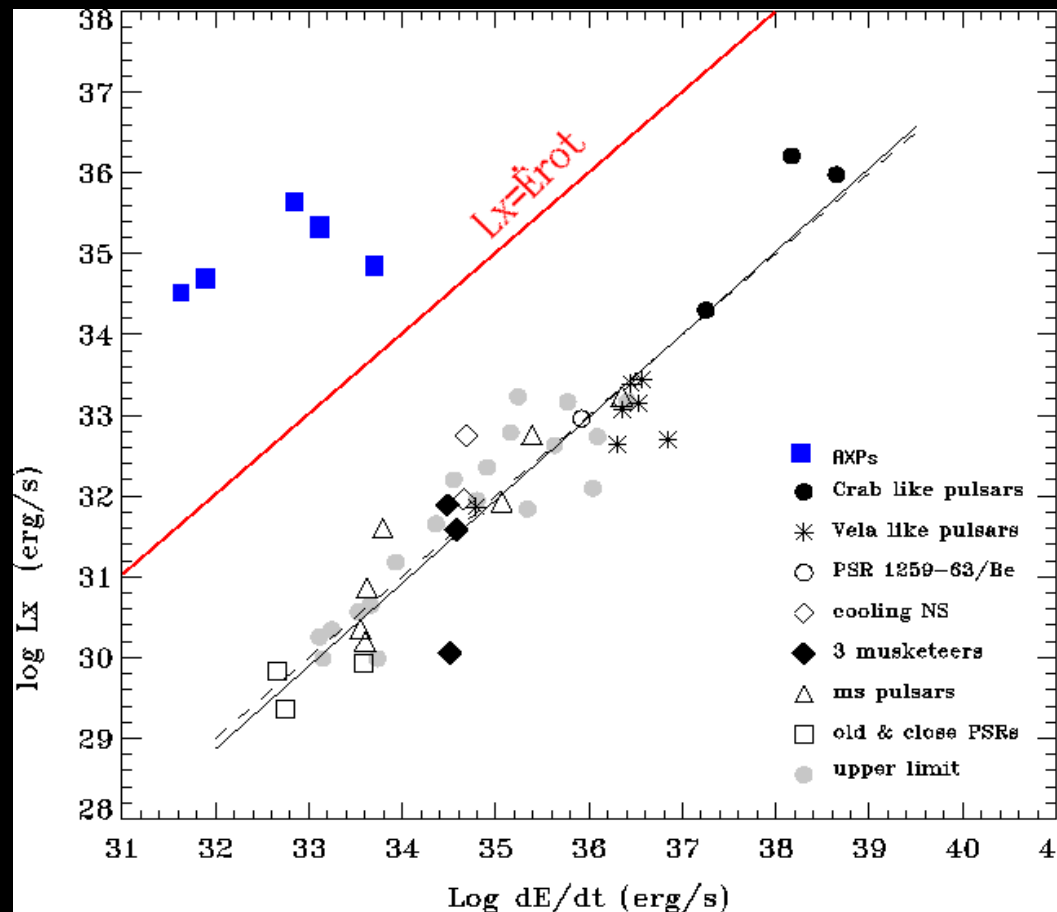
Recurrent flares reach $10^{41} \text{ erg/s} \sim 1000 L_{\text{Edd}}$, giant flares $10^{44} \text{ erg/s} \sim 10^6 L_{\text{Edd}}$

-> accretion energy ruled out



Persistent X-ray emission of AXPs

Rotational energy losses are orders of magnitudes smaller (10^{32-33} erg/s) than observed L_x !!



$$\dot{E}_{rot} = I\Omega\dot{\Omega} = 4\pi^2 I \dot{P} P^{-3}$$

Magnetars: Neutron Stars powered by magnetic energy

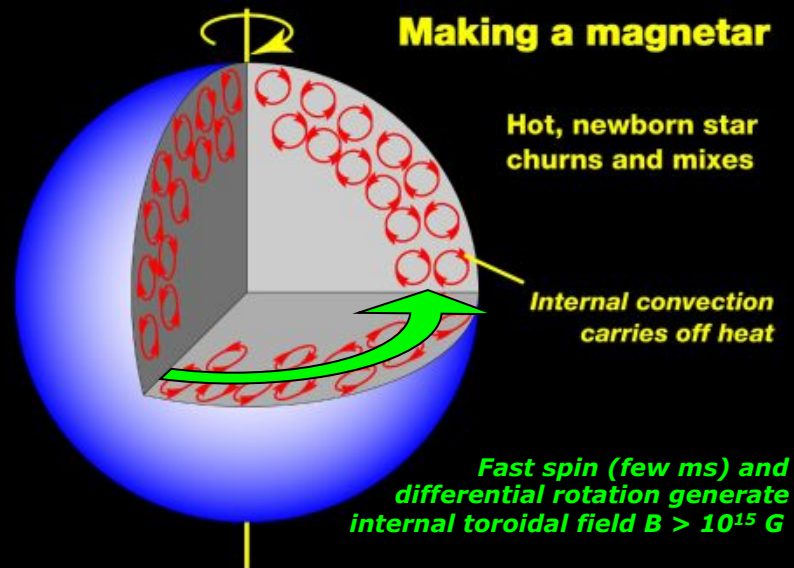
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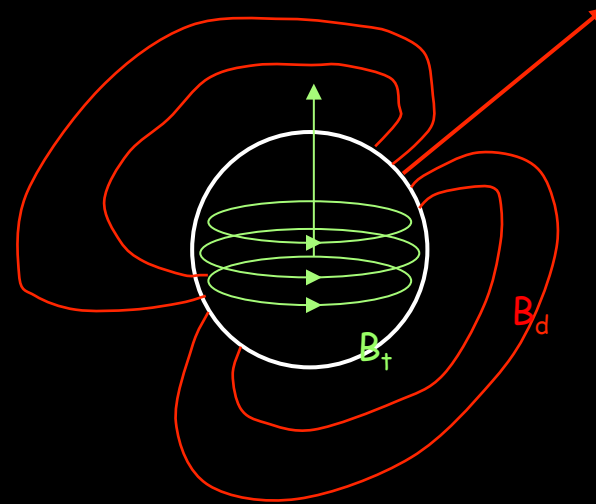
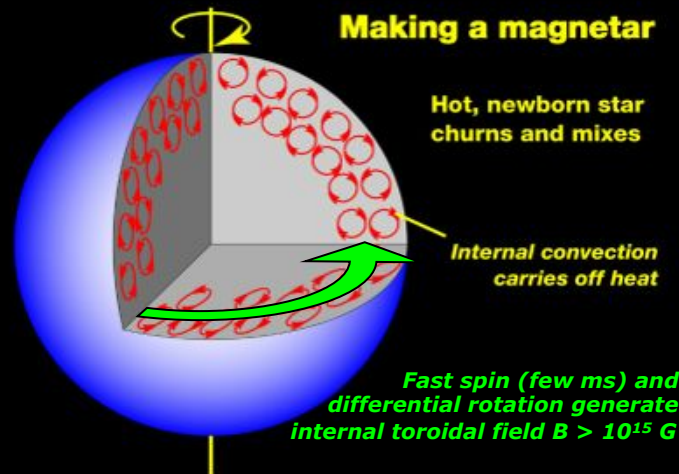


The B-field of Magnetars

Very strong internal B -fields in a newborn differentially rotating fast-spinning neutron star

For initial spin periods of $P_i \sim 1\text{--}2$ ms, differential rotation can store $\sim 10^{52} (P_i/1\text{ ms})^2$ ergs, that can be converted into a magnetic field of up to $3 \times 10^{17} (P_i/1\text{ ms})^{-1}$ G. (efficient dynamo might be limited to $\sim 3 \times 10^{16}$ G)

(Duncan & Thompson 1992)



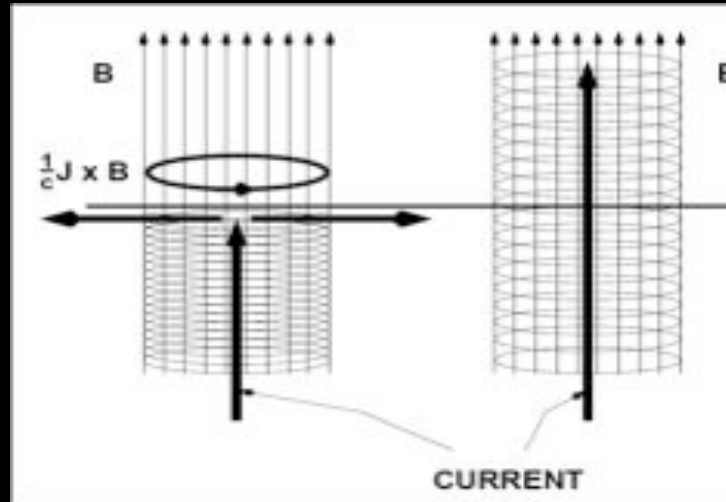
$B_d \sim 10^{14-15}$ G outer dipole field (spin-down, pulsations)
inferred from spin-down rate (and confirmed through the energetics and fast variability properties of the "ringing tail" of Giant Flares from SGRs)

$B_t > 10^{15}$ G inner toroidal field (energy reservoir):
lower limit from: $L(\text{persistent}) \times \text{age} \sim 10^{47}$ ergs

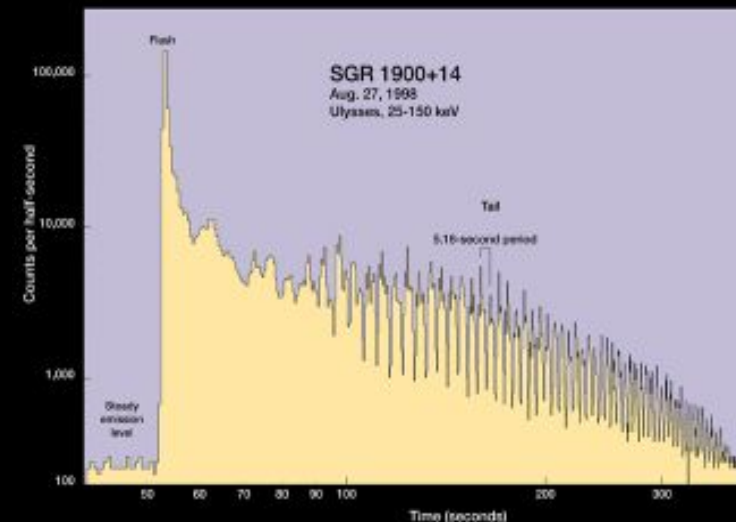
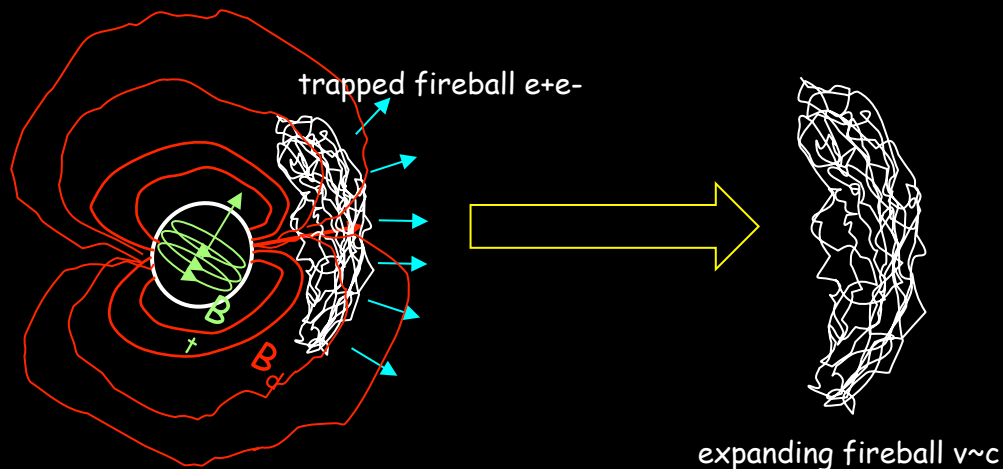
Bursts/Giant Flares in the Magnetar Model

Internal field evolves through ambipolar diffusion, building up magnetic stresses below NS crust

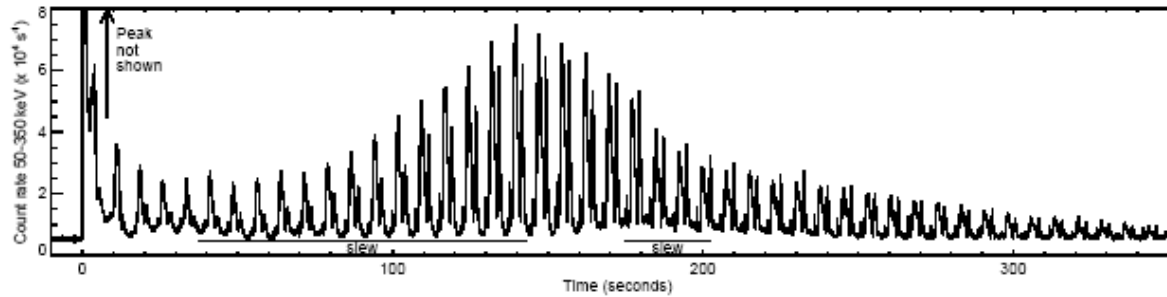
- Magnetic elicity fed to NS magnetosphere by local crustquakes of various amplitude, giving rise to **Recurrent bursts**



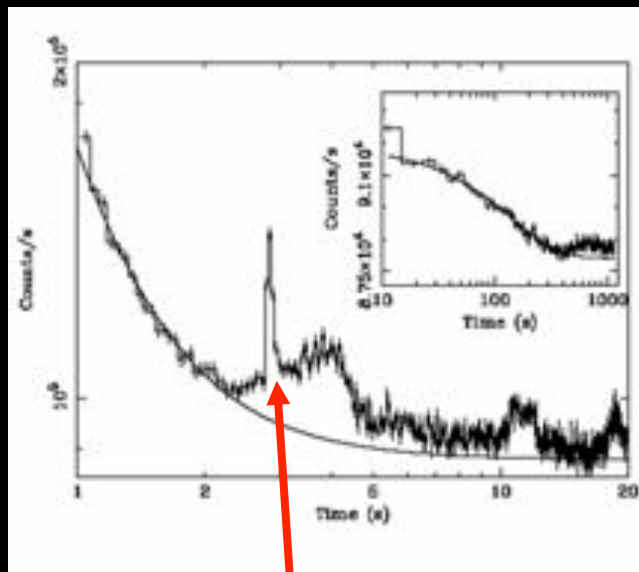
- **Giant flares** result from large-scale rearrangements of the core magnetic field:
- crustal yield strain is overcome, large-scale fracturing occurs that leads to sudden release of very large amounts of energy into the magnetosphere
- **fireball** of pair-dominated plasma expanding at $\sim c \rightarrow$ **Initial subsecond Spike** of Giant Flares
- **trapped fireball** gives rise to ~ 300 s oscillating tail



SGR 1806-20: Giant Flare of 2004 Dec 27



(Palmer et al. 2005
Hurley et al. 2005)



Moon reverberation seen !

Giant Flare Source	March 5, 1979 SGR 0526-66	August 27, 1998 SGR 1900+14	December 27, 2004 SGR 1806-20
Assumed distance	55 kpc	10 kpc	15 kpc
Initial Spike			
Duration (s)	~0.25	~0.35	~0.5
Peak luminosity (erg s ⁻¹)	$2.6 \cdot 10^{44}$	$> 3.7 \cdot 10^{44}$	$(2 \div 5) \cdot 10^{47}$
Fluence (erg cm ⁻²)	$4.5 \cdot 10^{-4}$	$> 5.5 \cdot 10^{-3}$	$0.6 \div 2$
Isotropic Energy (erg)	$1.6 \cdot 10^{44}$	$> 6.8 \cdot 10^{46}$	$(1.6 \div 5) \cdot 10^{48}$
Pulsating tail			
Duration (s)	~200	~200	~200
Fluence (erg cm ⁻²)	$1 \cdot 10^{-3}$	$4.2 \cdot 10^{-3}$	$5 \cdot 10^{-3}$
Isotropic Energy (erg)	$3.6 \cdot 10^{44}$	$5.2 \cdot 10^{46}$	$1.2 \cdot 10^{48}$
Spectrum	kT ~ 30 keV	kT ~ 20 keV	kT ~ 15-30 keV
Pulse Period (s)	8.1	5.15	7.56

(Mereghetti et al. 2005)